

Mu2e : Searching for $\mu^- N \rightarrow e^- N$ at Fermilab

Presented by Sophie Charlotte Middleton
Research Associate at Caltech
University of Michigan
October 2021

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Outline

Caltech



CHARGED LEPTON FLAVOR VIOLATION (CLFV)





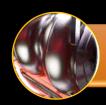
THE GLOBAL SEARCH FOR CLFV AND COMPLEMENTARITY



HOW CAN MU2E HELP US UNDERSTAND BEYOND STANDARD MODEL PHYSICS?



THE MU2E EXPERIMENT



THE MU2E-II EXPERIMENT



SUMMARY

smidd@caltech.edu

Mu2e: Searching for

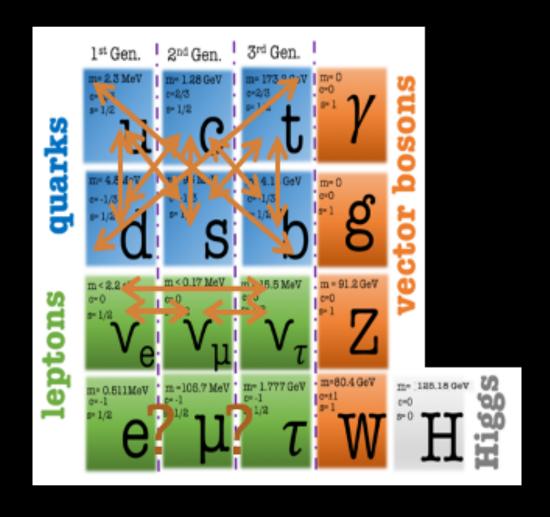
26 October 2021

Quark Flavor Violation



The quarks commit Flavour Violation -

Mixing strengths are parameterized by Cabibbo-Kobayashi-Maskawa (CKM) matrix:





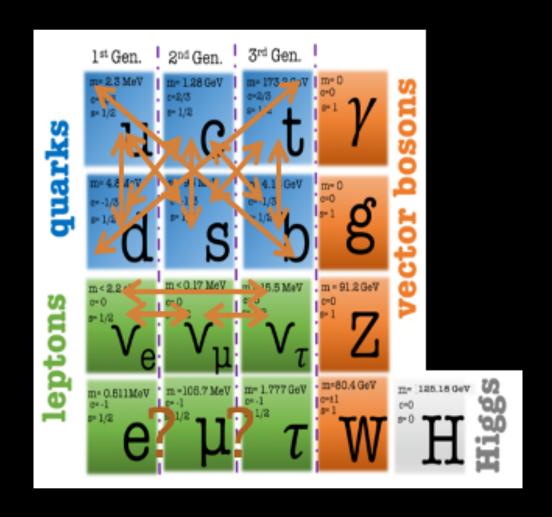
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$$(d',s',b') = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$





Quark Flavor Violation

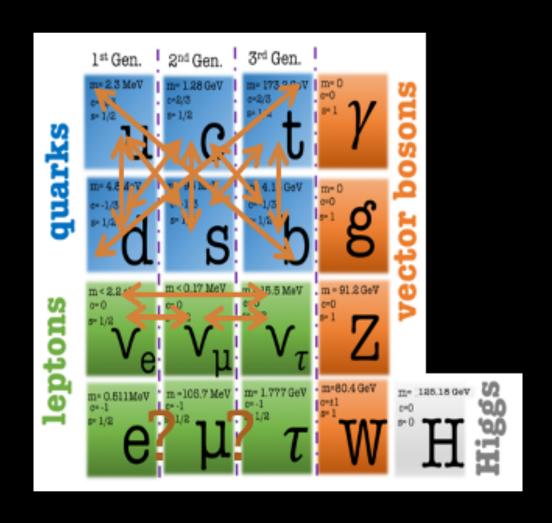


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→ almost diagonal.



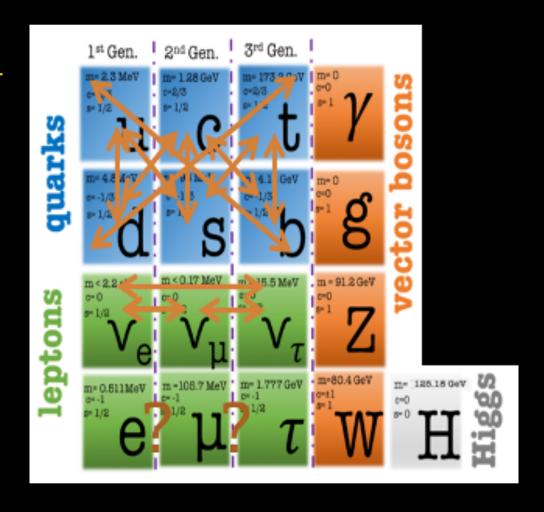


Neutral Lepton Flavor Violations



• ν oscillations \rightarrow Lepton Flavour Violation (LFV)

Mixing strengths parameterised by the Pontecorvo–Maki–Nakagawa–Sakata matrix (PMNS) matrix:





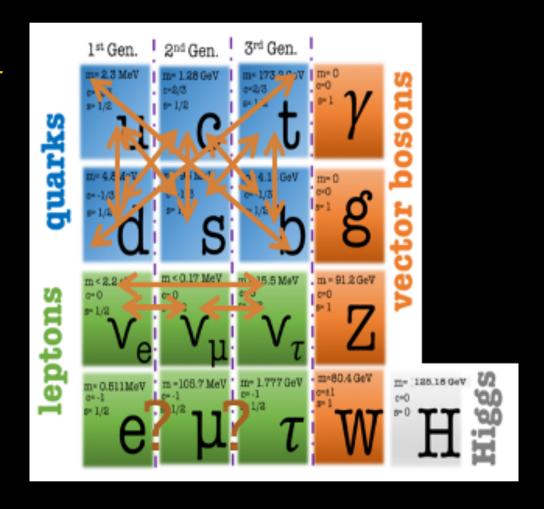
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Neutral Lepton Flavor Violations

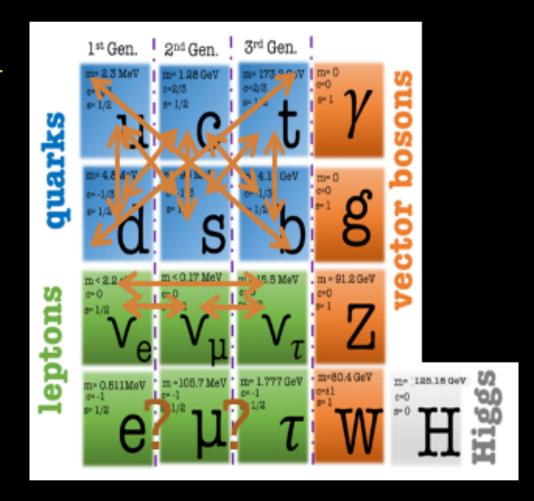


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→ The current best fit values imply that there is much more neutrino mixing than there is mixing between the quark flavours in the CKM matrix.





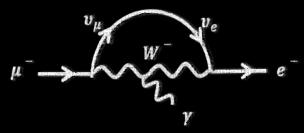


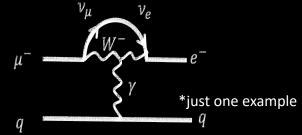
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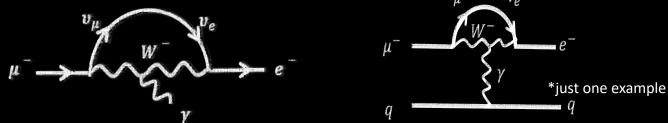








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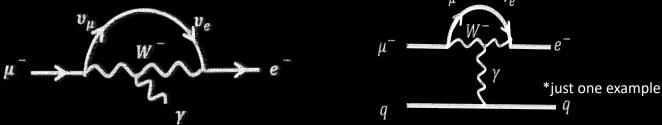


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$$B(\mu \to e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_{i=2,3} U_{\mu i}^* U_{ei} \frac{\Delta m_{1i}^2}{M_W^2} \right|^2$$

$$B(\mu \to e\gamma) = \frac{3\alpha}{32\pi} \left(\frac{1}{4}\right) \sin^2 2\theta_{13} \sin^2 \theta_{23} \left| \frac{\Delta m_{13}^2}{M_W^2} \right|^2$$

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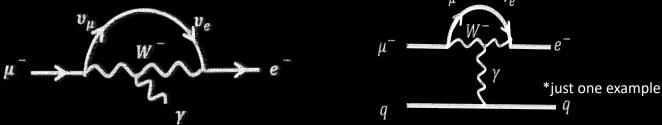
[1-4]

using best-fit values for neutrino data (m_{vj} for the neutrino mass and U_{ij} for the element of the PMNS matrix).





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If observed at Mu2e or Mu2e-II \rightarrow this would be an unambiguous sign of physics beyond the Standard Model (BSM).



BSM Scenarios

Nice overview: Lorenzo Calibbi, Giovanni Signorelli arXiv:1709.00294 (2018)



There are many well-motivated BSM theories which invoke CLFV mediated by (pseudo) scalar, (axial) vector, or tensor currents at rates close to current experimental limits i.e. $B \approx 10^{-15} - 10^{-17}$:



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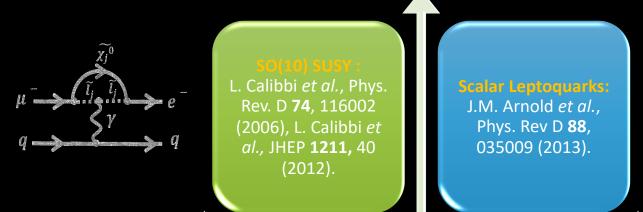


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A few examples – not an exclusive list!

Important point:

Mu2e is an indirect search for new physics, with sensitivity to lots of models. Even a null result would have huge implications!

Different neutrino mass-generating Lagrangians lead to very different rates for CLFV:

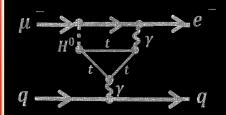
Nuclear Physics B (Proc. Suppl.) 248–250 (2014) 13–19

Extended Higgs/Gauge sector:

Left-Right Symmetric Models

C.-H. Lee *et al.*, Phys. ReV D **88**, 093010 (2013).

Littlest Higgs Blanke et al Phys.Polon.B41:657, 2010, arXiv:0906.5454v2 [hep-ph]



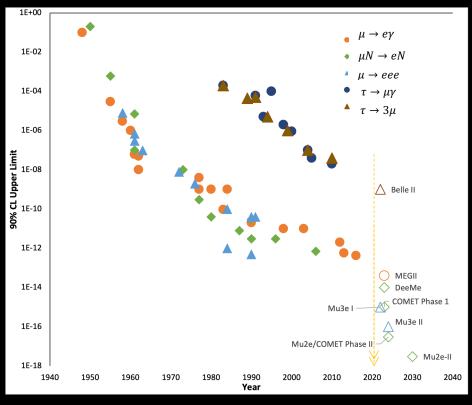


Mu2e: Searching for μ N→e N at milab - Sophie Middleton - smidd@caltech.edu



- $\mu^- N \rightarrow e^- N$ searches are crucial part of global program searching for CLFV.
- Muons offer more powerful probe for CLFV compared to taus.
- To elucidate the mechanism responsible for any CLFV must look at relative rates (if any) in different muon channels.

Mode	Current Limit (at 90% CL)	Future Proposed Limit	Future Experiment/s
$\mu^{\pm} o e^{\pm} \gamma$	4.2 x 10 ^{-13 [5]}	4 x 10 ⁻¹⁴	MEG II [8]
$\mu^- N \rightarrow e^- N$	7 x 10 ^{-13 [6]}	10 ⁻¹⁵ 10 ⁻¹⁷ 10 ⁻¹⁸	COMET Phase-I Mu2e [10] & COMET Phase-II [9] Mu2e-II
$\mu^+ \rightarrow e^+ e^+ e^-$	~10 ^{-12 [7]}	10 ⁻¹⁵ ~ 10 ⁻¹⁶	Mu3e



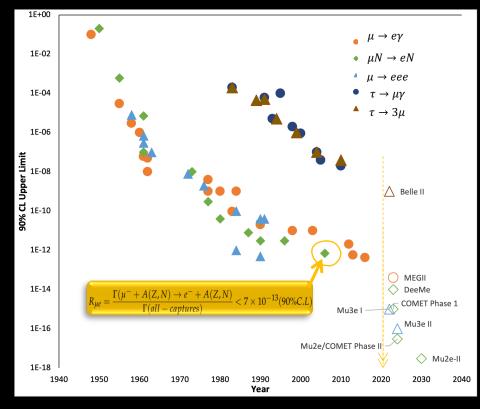
• Muon-to-electron sector provides powerful probes and complements collider searches for $\tau \to e \gamma$ or $\mu \gamma$ and $H \to e \tau$, $\mu \tau$, or μe .





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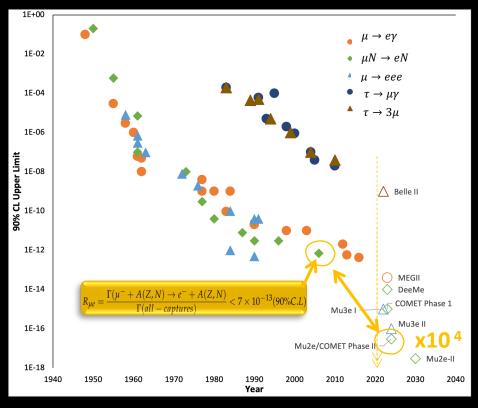
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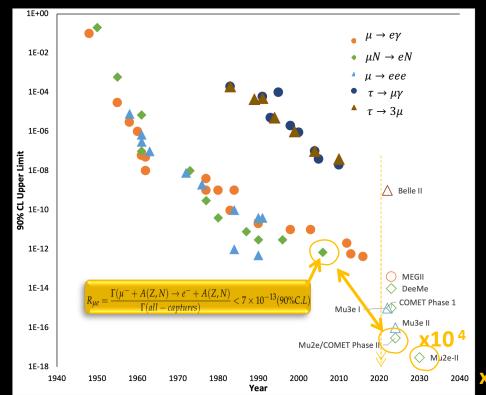
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Mu2e-II will use PIP-II, Snowmass Proposal in

• Muon-to-electron sector provides powerful probes and complements collider searches for $\tau \to e \gamma$ or $\mu \gamma$ and $H \to e \tau$, $\mu \tau$, or μe . progress!



Aside: Tau CLFV Searches

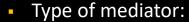
Taus Past/Present/Future: https://indico.fnal.gov/event/44457/?print=1



• Less stringent limits in 3rd generation, but here BSM effects may be higher.

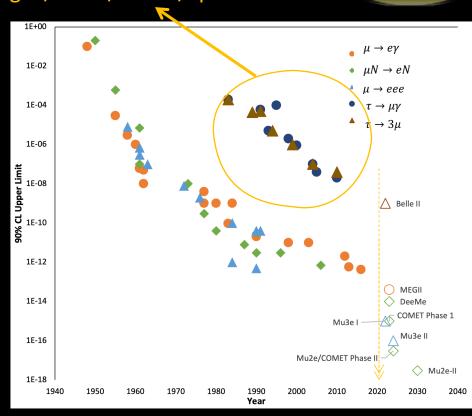
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$ au ightarrow \mu \mu \mu$	Belle – 2.1 x 10 ⁻⁸ BaBar – 3.3 x 10 ⁻⁸ LHCb – 4.6 x 10 ⁻⁸	< 1 x 10 ⁻⁹ (arXiv:1011.0352 arXiv:1808.10567)	Belle II, LHCb HL-LHC TauFV?
$ au ightarrow \mu \gamma$	4.4 x 10 ⁻⁸	~1 x 10 ⁻⁹	Belle II
$ au ightarrow e \gamma$	3.3 x 10 ⁻⁸	~3 x 10 ⁻⁹	Belle II

• τ LFV searches at Belle II will be extremely clean, with very little background (if any), thanks to pair production and double-tag analysis technique.

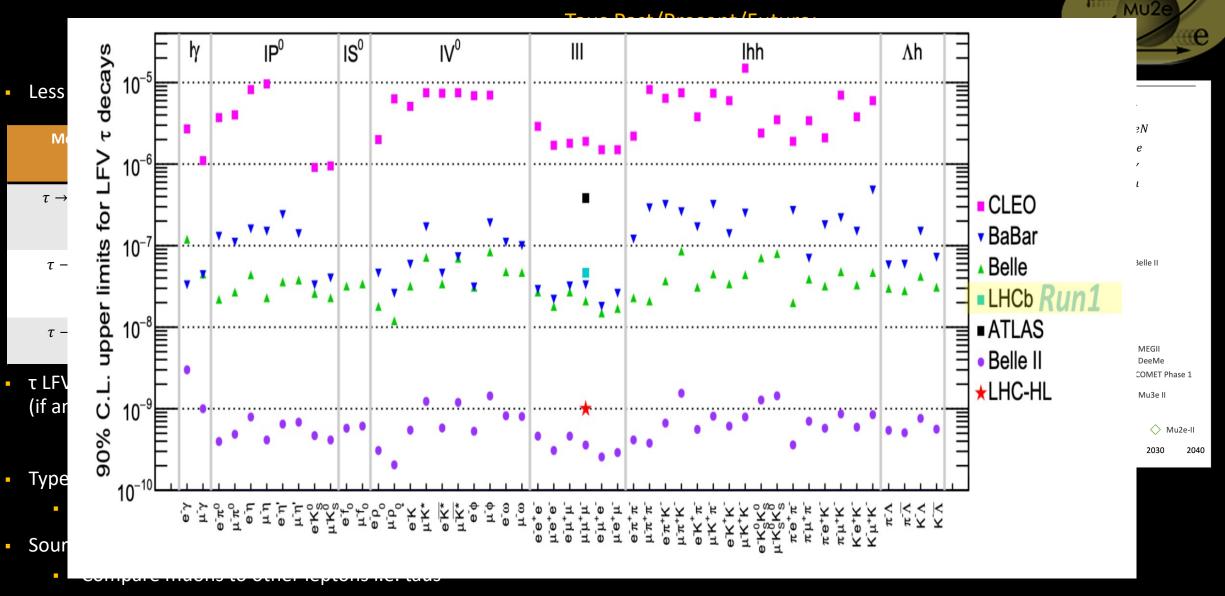


- Compare muon channels to each other
- Source of flavor violation:
 - Compare muons to other leptons i.e. taus





Aside: Tau CLFV Searches







A. de Gouvêa, P. Vogel arXiv:1303.4097

 For the purposes of discussion we can build a <u>Toy Lagrangian</u> which consists of 2 terms representing 2 types of physics process:

$$\mathcal{L}_{CLFV} = \frac{m_{\mu}}{(1+\kappa)\Lambda^2} \bar{\mu}_R \sigma_{\mu\nu} e_L F^{\mu\nu} + \frac{\kappa}{(1+\kappa)\Lambda^2} \bar{\mu}_L \gamma_{\mu} e_L (\sum_{q=u,d} \bar{q}_L \gamma_{\mu} q_L)$$





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$$q \longrightarrow q$$





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$$\chi^$$



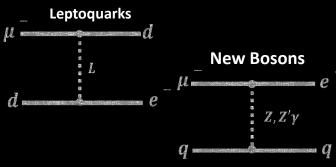


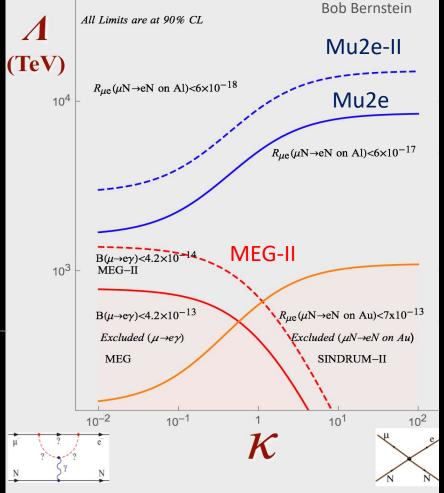
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 "Contact" i.e. 4 fermion terms Only $\mu^+\!\!\to e^+e^+e^-$ And $\mu^- N \to e^- N$





Mu2e e

Taken from: arXiv:0909.1333[hep-ph]

Ī		AC	RVV2	AKM	$\delta ext{LL}$	FBMSSM	LHT	RS	
	$D^0 - \bar{D}^0$	***	*	*	*	*	***	?	
	ϵ_K	*	***	***	*	*	**	***	
	$S_{\psi\phi}$	***	***	***	*	*	***	***	
	$S_{\phi K_S}$	***	**	*	***	***	*	?	
	$A_{\mathrm{CP}}\left(B o X_s\gamma ight)$	*	*	*	***	***	*	?	
	$A_{7,8}(B\to K^*\mu^+\mu^-)$	*	*	*	***	***	**	?	
	$A_9(B\to K^*\mu^+\mu^-)$	*	*	*	*	*	*	?	
	$B\to K^{(*)}\nu\bar\nu$	*	*	*	*	*	*	*	
	$B_s \to \mu^+ \mu^-$	***	***	***	***	***	*	*	
	$K^+ \to \pi^+ \nu \bar{\nu}$	*	*	*	*	*	***	***	
	$K_L o \pi^0 \nu \bar{\nu}$	*	*	*	*	*	***	***	
	$\mu \to e \gamma$	***	***	***	***	***	***	***	
	$ au ightarrow \mu \gamma$	***	***	*	***	***	***	***	
	$\mu + N \rightarrow e + N$	***	***	***	***	***	***	***	

Table 8: "DNA" of flavour physics effects for the most interesting observables in a selection of SUSY and non-SUSY models $\bigstar \star \star \star$ signals large effects, $\star \star$ visible but small effects and \star implies that the given model does not predict sizable effects in that observable.

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	$S_{\phi K_S}$	***	**	*	***	***	*	?	
	$A_{\mathrm{CP}}\left(B o X_s \gamma\right)$	*	*	*	***	***	*	?	
	$A_{7,8}(B\to K^*\mu^+\mu^-)$	*	*	*	***	***	**	?	
	$A_9(B\to K^*\mu^+\mu^-)$	*	*	*	*	*	*	?	
	$B \to K^{(*)} \nu \bar{\nu}$	*	*	*	*	*	*	*	
	$B_s \to \mu^+ \mu^-$	***	***	***	***	***	*	*	
	$K^+ \to \pi^+ \nu \bar{\nu}$	*	*	*	*	*	***	***	
	$K_L o \pi^0 u \bar{ u}$	*	*	*	*	*	***	***	
	$\mu \to e \gamma$	***	***	***	***	***	***	***	
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	$\mu + N \rightarrow e + N$	***	***	***	***	***	***	***	

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Discovery sensitivity across the board. Relative Rates however will be model dependent.





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	ϵ_K	*	***	***	*	*	**	***
	$S_{\psi\phi}$	***	***	***	*	*	***	***
	$S_{\phi K_S}$	***	**	*	***	***	*	?
	$A_{\mathrm{CP}}\left(B o X_s \gamma\right)$	*	*	*	***	***	*	?
	$A_{7,8}(B\to K^*\mu^+\mu^-)$	*	*	*	***	***	**	?
	$A_9(B o K^*\mu^+\mu^-)$	*	*	*	*	*	*	?
	$B \to K^{(*)} \nu \bar{\nu}$	*	*	*	*	*	*	*
	$B_s o \mu^+ \mu^-$	***	***	***	***	***	*	*
	$K^+ \to \pi^+ \nu \bar{\nu}$	*	*	*	*	*	***	***
	$K_L o \pi^0 u \bar{ u}$	*	*	*	*	*	***	***
	$\mu \to e \gamma$	***	***	***	***	***	***	***
	$ au ightarrow \mu \gamma$	***	***	*	***	***	***	***
	$\mu + N \rightarrow e + N$	***	***	***	***	***	***	***

Table 8: "DNA" of flavour physics effects for the most interesting observables in a selection of SUSY and non-SUSY models $\bigstar \star \star \star$ signals large effects, $\star \star$ visible but small effects and \star implies that the given model does not predict sizable effects in that observable.

the given model does not predict sizable effects in that observable

Discovery sensitivity across the board. Relative Rates however will be model dependent.

Model	$\mu \to eee$	$\mu N \to eN$	$\frac{\text{BR}(\mu \rightarrow eee)}{\text{BR}(\mu \rightarrow e\gamma)}$	$\frac{\operatorname{CR}(\mu N \to eN)}{\operatorname{BR}(\mu \to e\gamma)}$	arXiv
MSSM	Loop	Loop	$\approx 6 \times 10^{-3}$	$\frac{10^{-3}-10^{-2}}{10^{-3}-10^{-2}}$:17
Type-I seesaw	$Loop^*$	Loop*	$3 \times 10^{-3} - 0.3$	0.1 - 10	09.0
Type-II seesaw	Tree	Loop	$(0.1-3) \times 10^3$ $\mathcal{O}(10^-)$.00294v
Type-III seesaw	Tree	Tree	$\approx 10^3$	$\mathcal{O}(10^3)$	N
LFV Higgs	$\operatorname{Loop}^\dagger$	Loop*†	$\approx 10^{-2}$	$\mathcal{O}(0.1)$	[he
Composite Higgs	$Loop^*$	$Loop^*$	0.05-0.5	2-20	p-ph
	fr	om L. Calibbi and G	. Signorelli, Riv. Nuovo Cim	ento, 41 (2018) 71	



ми2е е

Taken from: arXiv:0909.1333[hep-ph]

	AC	RVV2	AKM	$\delta ext{LL}$	FBMSSM	LHT	RS
$D^0 - \bar{D}^0$	***	*	*	*	*	***	?
ϵ_K	*	***	***	*	*	**	***
$S_{\psi\phi}$	***	***	***	*	*	***	***
$S_{\phi K_S}$	***	**	*	***	***	*	?
$A_{\mathrm{CP}}\left(B o X_s \gamma\right)$	*	*	*	***	***	*	?
$A_{7,8}(B o K^*\mu^+\mu^-)$	*	*	*	***	***	**	?
$A_9(B o K^*\mu^+\mu^-)$	*	*	*	*	*	*	?
$B \to K^{(*)} \nu \bar{\nu}$	*	*	*	*	*	*	*
$B_s o \mu^+ \mu^-$	***	***	***	***	***	*	*
$K^+ \to \pi^+ \nu \bar{\nu}$	*	*	*	*	*	***	***
$K_L o \pi^0 u \bar{ u}$	*	*	*	*	*	***	***
$\mu \to e \gamma$	***	***	***	***	***	***	***
$ au ightarrow \mu \gamma$	***	***	*	***	***	***	***
$\mu + N \rightarrow e + N$	***	***	***	***	***	***	***

Table 8: "DNA" of flavour physics effects for the most interesting observables in a selection of SUSY and non-SUSY models $\bigstar \star \star \star$ signals large effects, $\star \star$ visible but small effects and \star implies that the given model does not predict sizable effects in that observable.

the given model does not predict sizable effects in that observable

Discovery sensitivity across the board. Relative Rates however will be model dependent.

Theory on complementarity -> http://arxiv.org/abs/2010.00317v2

Model	$\mu \to eee$	$\mu N \to e N$	$\frac{\text{BR}(\mu{\rightarrow}eee)}{\text{BR}(\mu{\rightarrow}e\gamma)}$	$\frac{\operatorname{CR}(\mu N \to eN)}{\operatorname{BR}(\mu \to e\gamma)}$
MSSM	Loop	Loop	$\approx 6 \times 10^{-3}$	$10^{-3} - 10^{-2}$
Type-I seesaw	$Loop^*$	Loop^*	$3 \times 10^{-3} - 0.3$	0.1-10
Type-II seesaw	Tree	Loop	$(0.1-3)\times10^3$	$\mathcal{O}(10^{-2})$
Type-III seesaw	Tree	Tree	$\approx 10^3$	$O(10^{3})$
LFV Higgs	$\operatorname{Loop}^\dagger$	Loop* †	$\approx 10^{-2}$	$\mathcal{O}(0.1)$
Composite Higgs	$Loop^*$	Loop*	0.05 - 0.5	2-20
		from L. Calibbi and G.	ignorelli, Riv. Nuovo Cin	nento, 41 (2018) 71

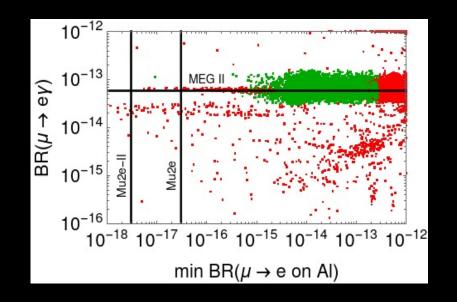


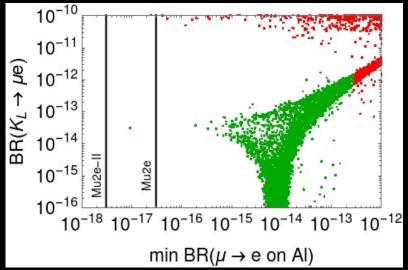
Example: Leptoquarks

Heeck & Teresi, 1808.07492



 Pati-Salam Leptoquarks: B-meson anomalies could be explained by two scalar leptoquarks, whose couplings enter neutrino masses as well. Type-II seesaw dominance is favored.





Plots: Julian Heeck

- Flavor structure fixed by neutrino mass/mixing; scale Λ fixed to explain B-meson anomaly R(K).
- Predicts testable rates in Mu2e!
- Explain B meson anomalies through adding 2 scalar leptoquarks: Bigaran, Gargalionis, Volkas, 1906.01870 \rightarrow B($\mu N \rightarrow eN$) < 3 x 10⁻¹³

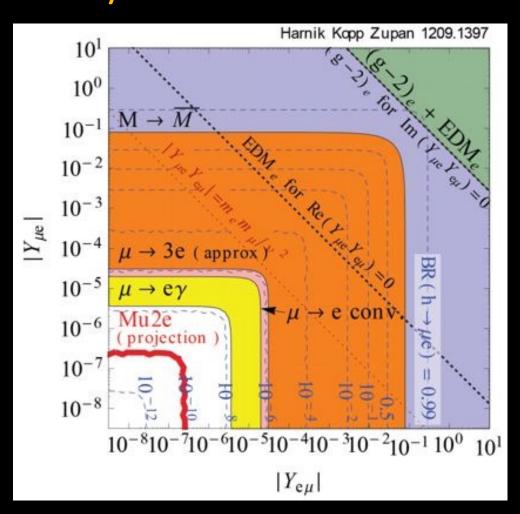


Example: Constraining Flavor Violating Higgs

Decays

Harnik, J. Kopp and J. Zupan, JHEP 3, 26 (2013).





- Can also help constrain flavor violating Higgs decays.
- Higgs LFV decays arise in many frameworks of New Physics at the electroweak scale such as two Higgs doublet models, extra dimensions, or models of compositeness.
- Current $\mu \rightarrow e$ conversion implies:

$$\sqrt{|Y_{\mu e}|^2 + |Y_{e\mu}|^2} < 4.6 \times 10^{-5}$$

• Mu2e is expected to be sensitive to:

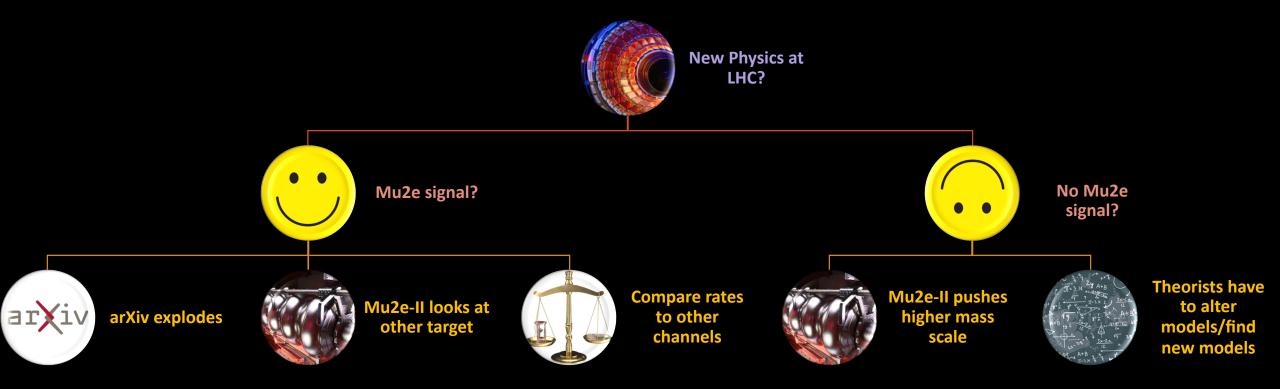
$$\sqrt{|Y_{\mu e}|^2 + |Y_{e\mu}|^2} > \text{few} \times 10^{-7}.$$

- Where $|Y_{\mu e}|$ and $|Y_{e\mu}|$ are flavor-violating Yukawa couplings for a 125 GeV Higgs boson i.e. $h \to \mu e$.
- Very strong limits on LFV Higgs decays for 1st-2nd generation



Possibilities







Complementarity in Target Materials

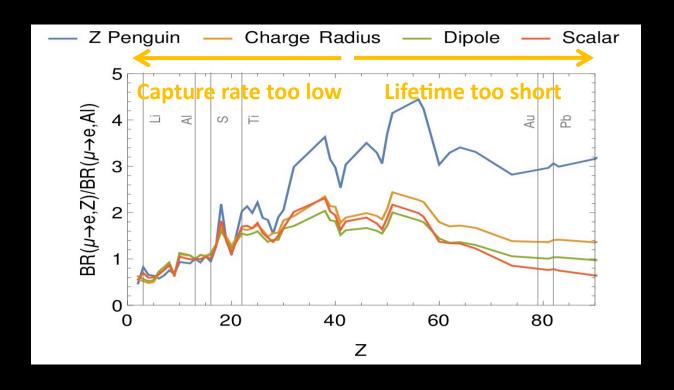
$$\mathsf{BR}(\mu \to \mathsf{e}) \propto |\mathsf{DC_{DL}} + \mathsf{S^pC_{S,L}^p} + \mathsf{V^pC_{V,R}^p} + \mathsf{S^nC_{S,L}^n} + \mathsf{V^nC_{V,R}^n}|^2 + (\mathsf{L} \leftrightarrow \mathsf{R})$$

	3	D	V ¹	V ²
$\frac{B(\mu \rightarrow e, Ti)}{B(\mu \rightarrow e, Al)}$	$1.70 \pm 0.005_{y}$	1.55	1.65	2.0
D(D()	$0.69 \pm 0.02_{\rho_n}$	1.04	1.41	$2.67 \pm 0.06_{\rho_n}$

If we do see a signal in Al at Mu2e, what can we do?:

- Various operator coefficients add coherently in the amplitude.
- Weighted by nucleus-dependent functions.
- → Requires measurements of conversion rate in other target materials!
- Need to choose a target which is sensitive to directions Al is "blind" to

This is something we need to think about for Mu2e-II – our extended program (see final slide).

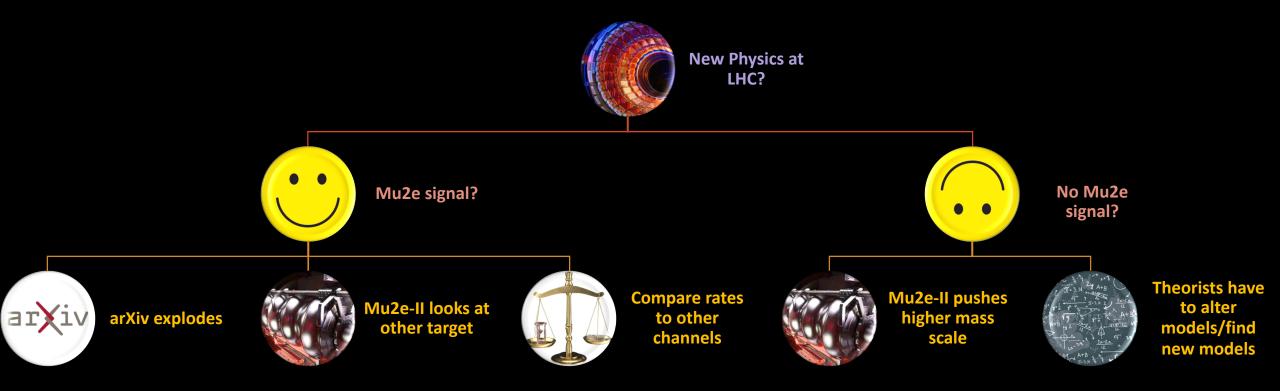


V. Cirigliano, S. Davidson, Y. Kuno, Phys. Lett. B 771 (2017) 242 S. Davidson, Y. Kuno, A. Saporta, Eur. Phys. J. C78 (2018) 109 Kitano et al 2002



Possibilities







Other Physics Searches at Mu2e



$$\mu^- N \rightarrow e^+ N'$$

$$R_{\mu^-e^+}^{\mathrm{Ti}} \equiv rac{\Gamma(\mu^- + \mathrm{Ti} o e^+ + \mathrm{Ca})}{\Gamma(\mu^- + \mathrm{Ti} o
u_\mu + \mathrm{Sc})} < \left\{ egin{array}{l} 1.7 imes 10^{-12} & \mathrm{(GS, 90\% \ CL)} \ 3.6 imes 10^{-11} & \mathrm{(GDR, 90\% \ CL)} \end{array}
ight.$$

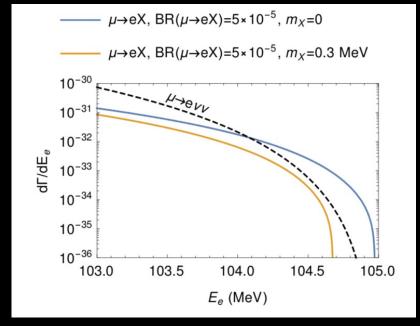
- This conversion violates both lepton number (ΔL = 2) and lepton flavor conservation, and can be mediated by Majorana neutrinos through a type-1 see saw mechanism or new particle at > TeV scale.
- The Mu2e sensitivity to μ \rightarrow e+ extends beyond the current best limit: Phys Rev Lett B 412 p 334-338 [13]

• < $m_{e\mu}>$ effective Majorana neutrino mass scale sensitivity down to the MeV region, surpassing the < $m_{\mu\mu}>$ sensitivity in the kaon sector which is limited to the GeV region

Yeo et al 2017 https://arxiv.org/abs/1705.07464

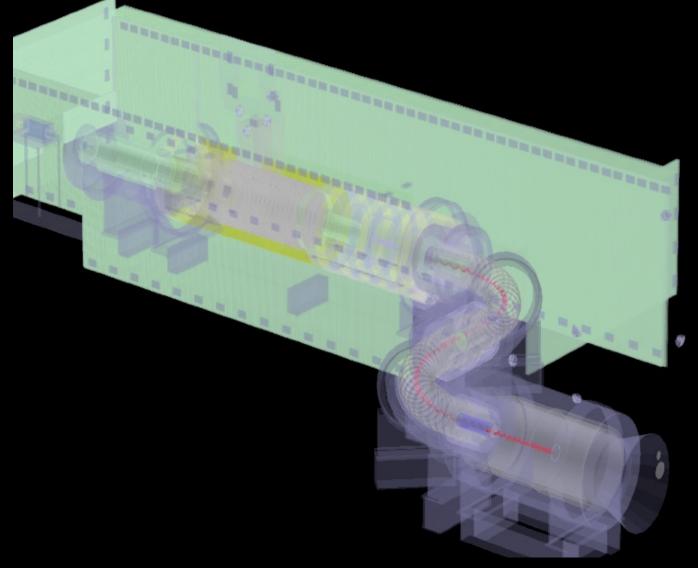
$$\mu^- N \rightarrow eXN$$

- Where X is a new light boson (or axion).
- Currently understanding feasibility
- Example parameterization: arXiv: 1110.2874









Experimental Strategy

How will Mu2e try to measure this process?



Muonic Atoms

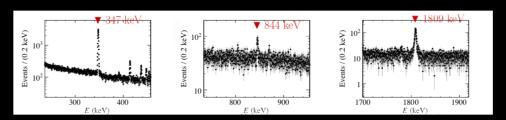


• The $\mu \rightarrow e$ conversion rate is measured as a ratio to the muon capture rate on the same nucleus:

$$R_{\mu e} = rac{\Gamma(\mu^- + A(Z,N)
ightarrow e^- + A(Z,N)}{\Gamma(all-captures)}$$

- Low momentum (-) muons are captured in the target atomic orbit and quickly (\sim fs) cascades to 1s state.
- Lifetime of muonic aluminium = 864 ns
- In aluminum:
 - 39 % Decay : $\mu + N \rightarrow e + \bar{v}_e + v_\mu$ (Background)
 - 61 % Capture : $\mu + N \rightarrow v_{\mu} + N'$ (Normalization)
 - The Signal : $\mu + N \rightarrow e + N$ (Conversion)

Normalization = from X -rays emitted when muon stops in Al.



Signal is monoenergetic electron consistent with:

$$E_e=m_\mu-E_{recoil}-E_{1S\,B.E}$$
 , e.g For Al: E_e = 104.97 MeV.

- Will be smeared by detector and stopping target effects.
- Nucleus coherently recoils off outgoing electron; it does not break-up!

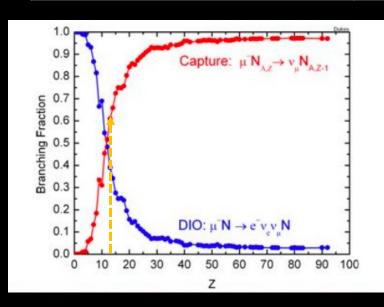


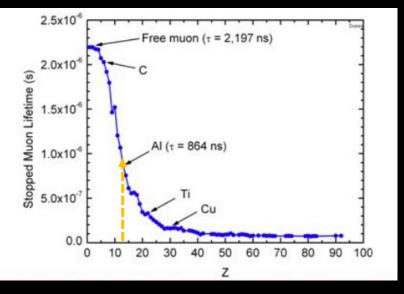


Why Aluminum?



Practical Advantages	Physics Advantages
Chemically Stable	Conversion energy such that only tiny fraction of photons produced by muon radiative capture.
Available in required size/shape/thickness	Muon lifetime long compared to transit time of prompt backgrounds.
Low cost	Conversion rate increases with atomic number, reaching maximum at Se and Sb, then drops. Lifetime of muonic atoms decreases with increasing atomic number.
	Lifetime of muonic atom sits in "goldilocks" region i.e. neither longer than 1700 ns pulse spacing and greater than our pionic live gate.





The lifetime of a muon in a muonic atom decreases with increasing atomic number.



Mu2e gets 8kW, 8GeV Protons from the Fermilab booster:

• Mu2e will acquire ϑ (10²⁰) Protons on Target to achieve design goal



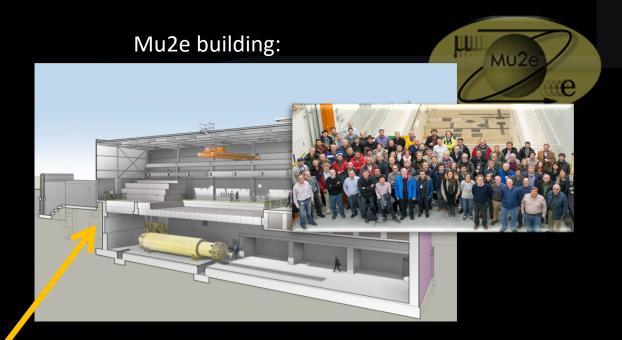




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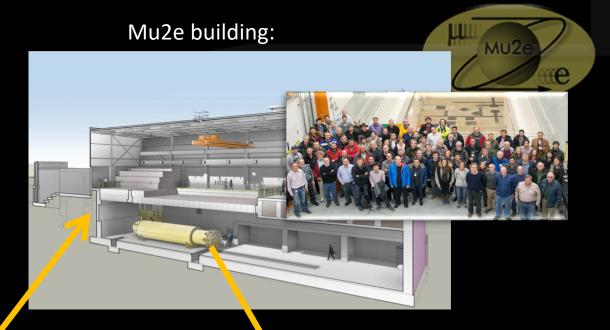




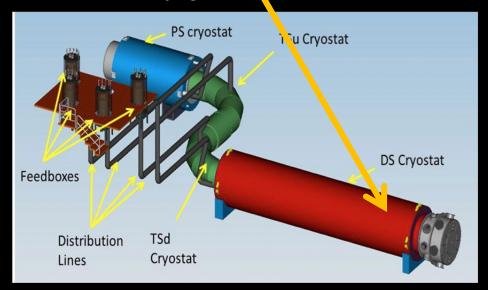
Mu2e gets 8kW, 8GeV Protons from the Fermilab booster:

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Mu2e:





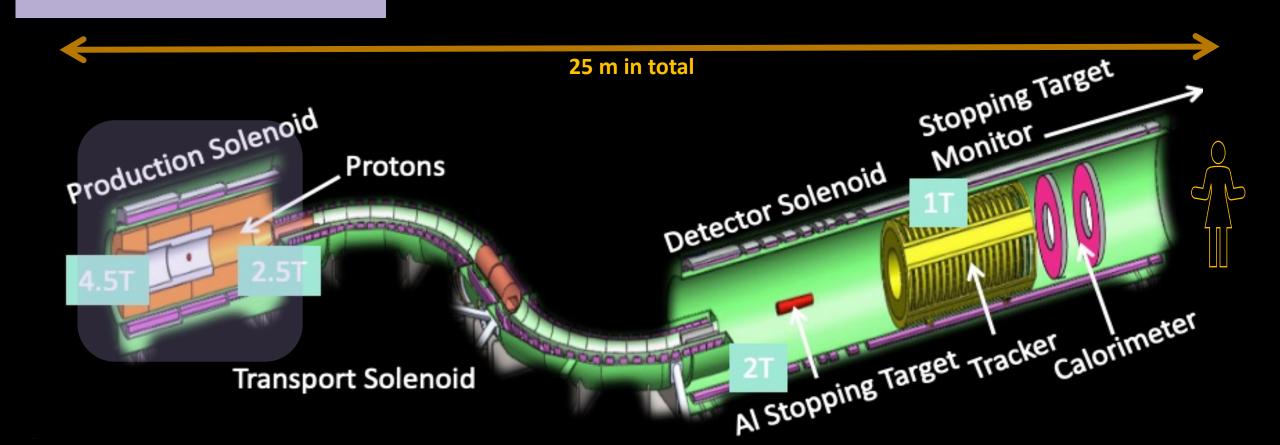
The Mu2e Experiment

V. Lobashev & R. Djilkibaev (Sov. J. Nucl. Phys. 49(2), 384 (1989))



Production Solenoid

- 8 GeV Protons enter, pions produced, decay to muons
- Graded magnetic field reflects pions/muons to transport solenoid



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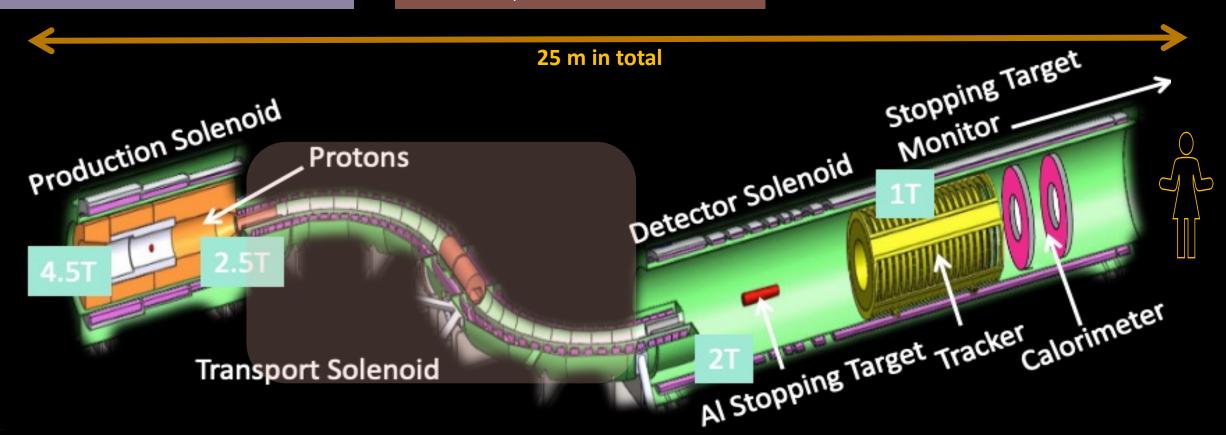


Production Solenoid

- 8 GeV Protons enter, pions produced, decay to muons
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Transport Solenoid:

- "S" shape removes line of sight backgrounds
- Windows remove anti-protons
- Collimators help select low momentum, negative muons and "focus" on detector solenoid aperture.



The Mu2e Experiment

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Production Solenoid:

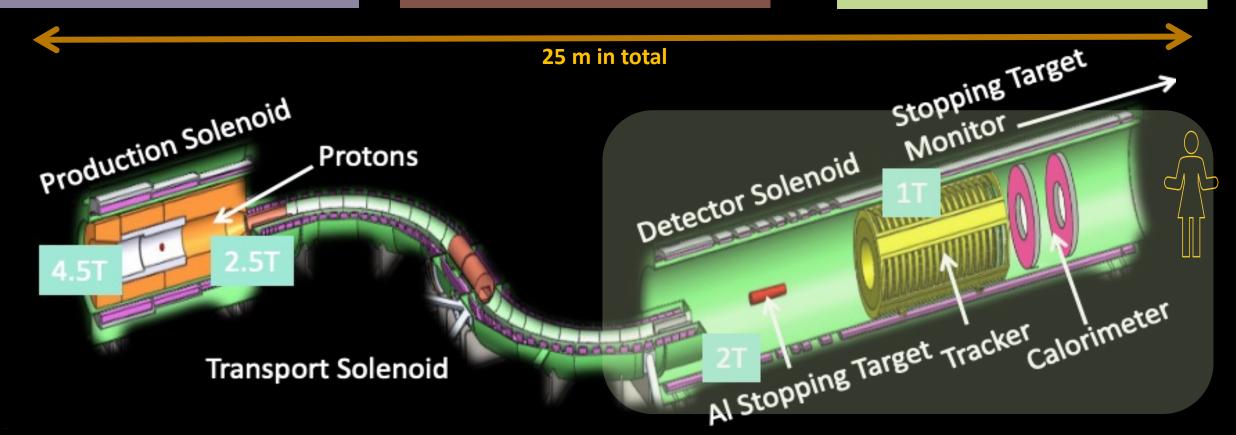
- 8 GeV Protons enter, pions produced, decay to muons
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Transport Solenoid:

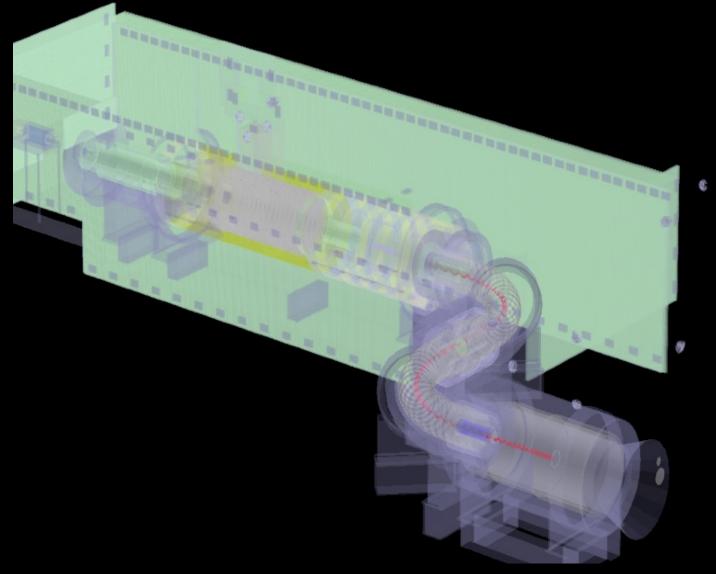
- "S" shape removes line of sight backgrounds
- Windows remove anti-protons
- Collimators help select low momentum, negative muons and "focus" on detector solenoid aperture.

Detector Solenoid:

- Al Stopping Target made of thin foils captures the muons
- Graded magnetic field "focusses" electrons on tracker
- Straw tracker and calorimeter measurementum







Beamline & Solenoids: Status

What is the current status of the beamline components?

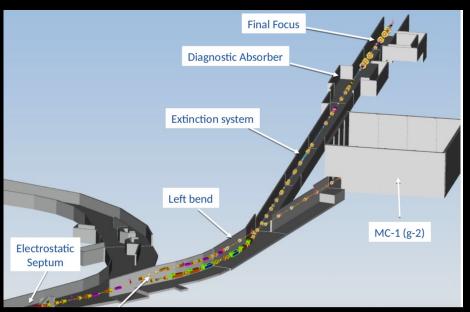


Beamline: Status





Quadrupole in Delivery Ring used for resonant extraction



The beamline installation is almost complete:

- Vacuum System installed
- Instrumentation upstream of diagnostic absorber in progress

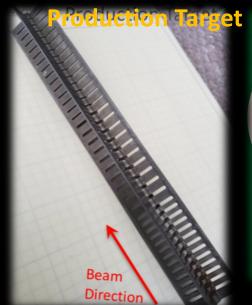


Production Target: Status

- Made of tungsten, completed in April 2021 10% of beam power into the target,
- Heats up to 1700 °C (~3100 F),
- Production Solenoid must be radiatively cooled,
- Average power density ~150 MW/m³

Production Target & Frame

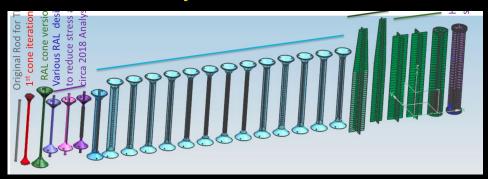






Production Target in support

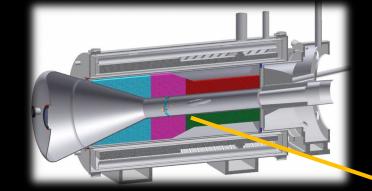
Many iterations



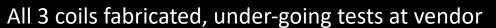
https://www.symmetrymagazine.org/article/a-robot-ballet



Production Solenoid: Status







Heat & Radiation Shield





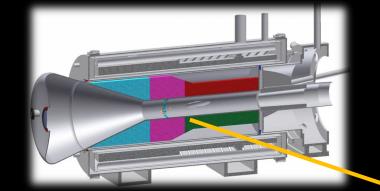
Production Target & Frame

https://www.symmetrymagazine.org/article/a-robot-ballet



Nice video of robotic extraction of target

Production Solenoid: Status





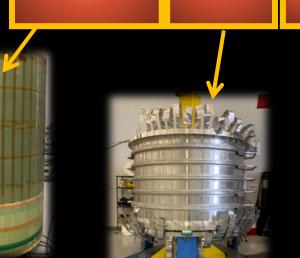


All 3 coils fabricated, under-going tests at vendor

Heat & Radiation Shield



Production Target & Frame







https://www.symmetrymagazine.org/article/a-robot-ballet

Caltech

Nice video of robotic extraction of target

Transport Solenoid: Status







- TSu and TSd cold-masses assembled.
- Testing almost complete
- Outer thermal shield will be split and reassembled around the TSu cold-mass alongside in image.







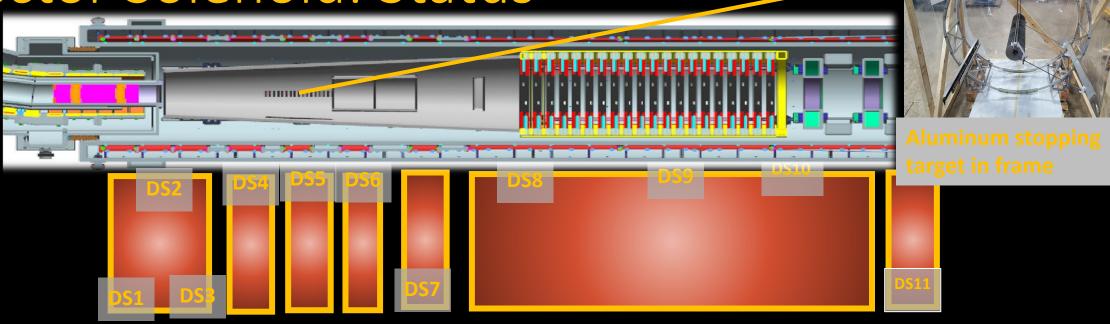
Nice video of the TS being lifted at FNAL in July:







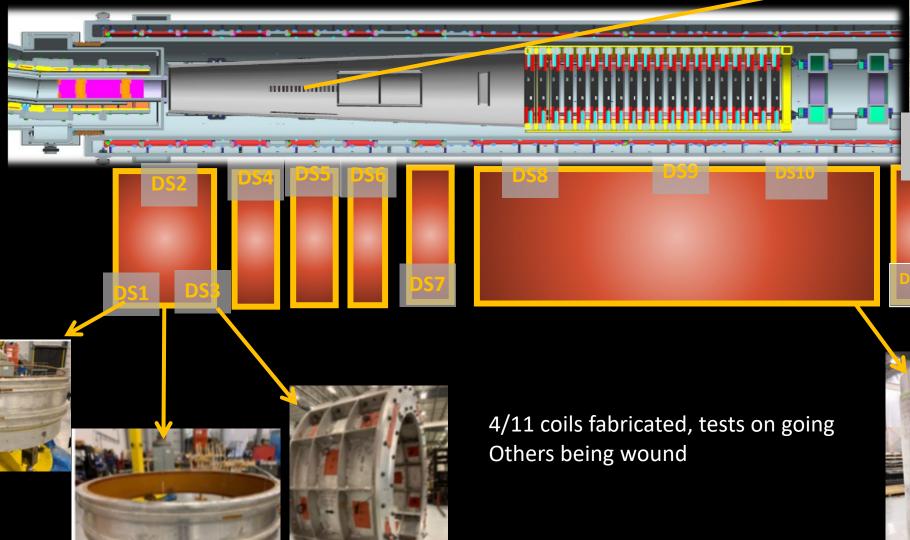
Detector Solenoid: Status





Detector Solenoid: Status

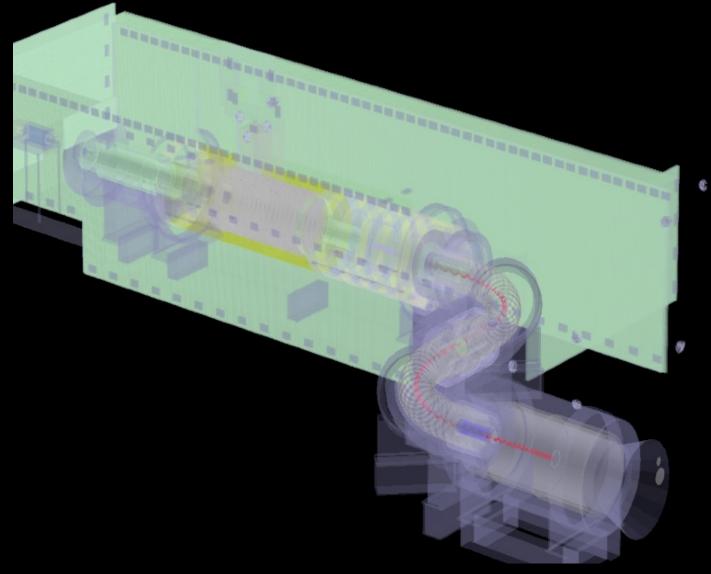
Caltecn



1511

for μ N→e N at Fermilab - Sophie Middleton smidd@caltech.edu





"Background free" design

How do we ensure Mu2e is "background free"? Why have we designed our detectors in this way?



Removing Backgrounds



Beam delivery and detector systems optimized for high intensity, pure muon beam – must be "background free":

Intrinsic:

 Scale with number of stopped muons.

- Late arriving :
 - Scale with number of late protons/ extinction performance

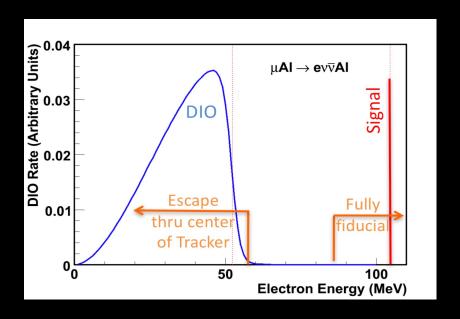
Туре	Source	Mitigation	Yield (over lifetime of experiment)
Intrinsic	Decay in Orbit (DIO)	Tracker Deign/ Resolution	0.144 \pm 0.028 (stat) \pm 0.11 (sys)
Late Arriving	Pion Capture	Beam Structure /Extinction	0.021 \pm 0.001 (stat) \pm 0.002 (sys)
	Pion Decay in Flight	-	$0.001 \pm < 0.001$
Other	Anti-proton	Thin Absorber Windows	0.04 \pm 0.022 (stat) \pm 0.020 (sys)
	Cosmic Rays	Active Veto System	0.209 \pm 0.0022 (stat) \pm 0.055 (sys)





In Aluminium 39% of stopped muons will decay in orbit (DIO):

 Free muon decay: peak electron energy far below our signal energy (peaks 52.8 MeV).





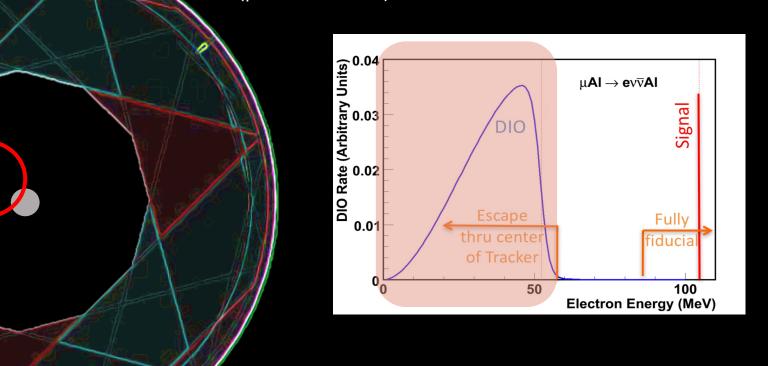


Transverse Plane

Michel Electron (52MeV/c)

In Aluminium 39% of stopped muons will decay in orbit (DIO):

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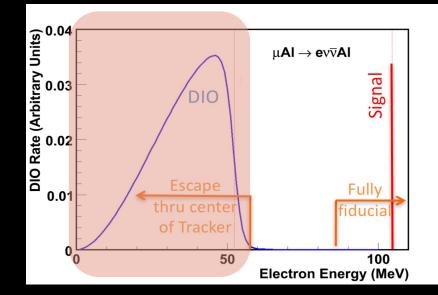


Transverse Plane

Michel Electron (52MeV/c)

In Aluminium 39% of stopped muons will decay in orbit (DIO):

Free muon decay: peak electron energy far below our signal energy (peaks 52.8 MeV).



Annular Design → Excludes low momentum electrons via hollow centre:

- Inner 38 cm un-instrumented.
- Reduces need to reject $\sim 10^{18}$ to $\sim 10^5$.
- Blind to > 99% of DIO spectrum.

Caltech

Mu2e: Searching for μ N→e N at Fermilab - Sophie Middleton - smidd@caltech.edu

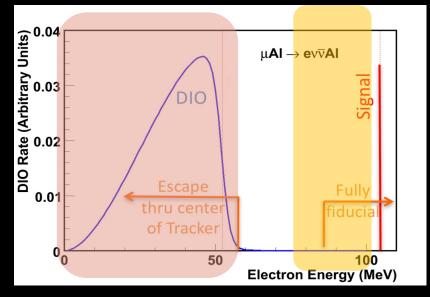




- Michel Electron (52MeV/c)
- Problematic Tail (>75MeV/c)

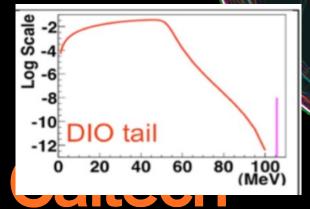
If muon bound in atomic orbit, the outgoing electron can exchange momentum with the nucleus.

Electron could have energy close to signal.



Annular Design → cannot fully exclude electrons in the recoil tail.

Recoil tail:



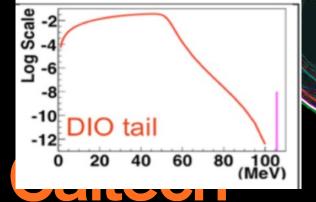




- Michel Electron (52MeV/c)
- Problematic Tail (>75MeV/c)

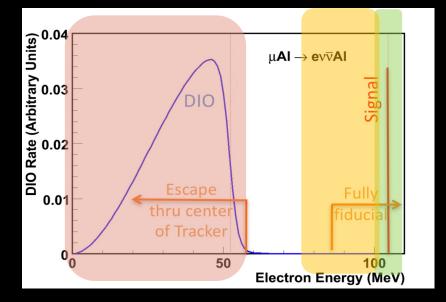
Signal (105MeV/c)

Recoil tail:



 If muon bound in atomic orbit, the outgoing electron can exchange momentum with the nucleus.

Electron could have energy close to signal.

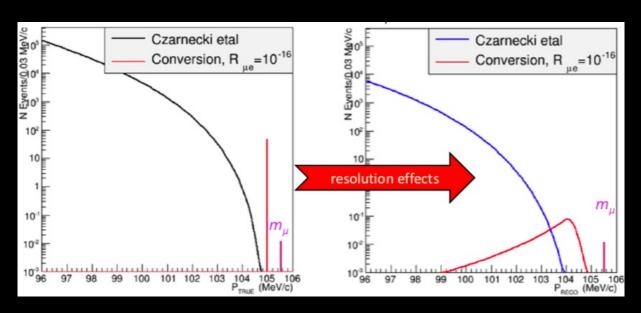


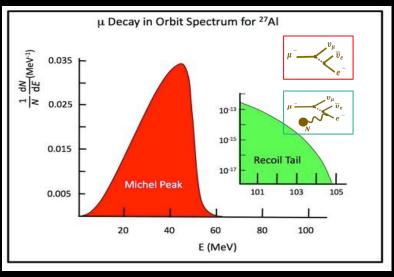
Annular Design → cannot fully exclude electrons in the recoil tail.



The differential energy spectrum of DIO electron spectrum has been parameterized in A. Czarnecki et al., "Muon decay in orbit: Spectrum of high-energy electrons," Phys. Rev. D 84 (Jul, 2011).

Necessitates tracker resolution of better than 200 KeV/c









■ To remove remaining DIOs momentum resolution < 200 KeV/c achieved by:





- To remove remaining DIOs momentum resolution < 200 KeV/c achieved by:
- Low Mass → Minimizes scattering and energy loss :
 - Entire Detector Solenoid held under vacuum ($\sim 10^{-4}$ torr).
 - Ultra low mass tracker.





- To remove remaining DIOs momentum resolution < 200 KeV/c achieved by:
- Low Mass → Minimizes scattering and energy loss :
 - Entire Detector Solenoid held under vacuum ($\sim 10^{-4}$ torr).
 - Ultra low mass tracker.
- 2. Segmented -> Handle high rates and provide high-precision momentum measurements.



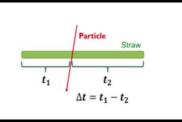
The Tracker: Design

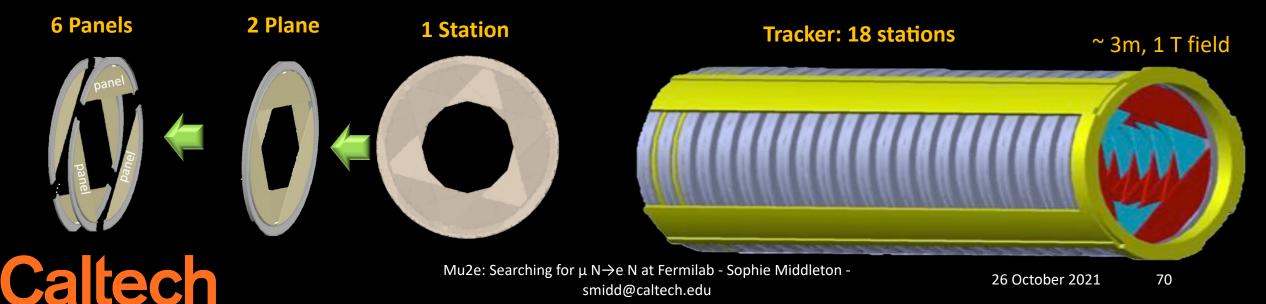
Mu2e e

- Tracker is constructed from self-supporting panels of low mass straws tubes detectors
- 18 stations, 2 planes per station, 6 panels per plane, 96 straws per panel.
- Straw drift tubes aligned transverse to the axis of the Detector Solenoid.
 - 1m, 5 mm diameter straw
 - Walls: 12 mm Mylar + 3 mm epoxy
 - 25 mm Au-plated W sense wire
 - 33 117 cm in length
 - 80:20 Ar:CO₂ with HV < 1500 V
 - Straw wall thickness of 15 μm has never been done before
- Charged particles ionize gas drift to wire detect signals!

The Straws:









- To remove remaining DIOs momentum resolution < 200 KeV/c achieved by:
- Low Mass → Minimizes scattering and energy loss :
 - Entire Detector Solenoid held under vacuum ($\sim 10^{-4}$ torr).
 - Ultra low mass tracker.
- 2. Segmented -> Handle high rates and provide high-precision momentum measurements.

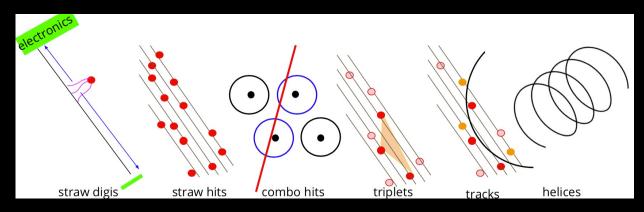




- To remove remaining DIOs momentum resolution < 200 KeV/c achieved by:
- Low Mass → Minimizes scattering and energy loss :
 - Entire Detector Solenoid held under vacuum ($\sim 10^{-4}$ torr).
 - Ultra low mass tracker.
- 2. Segmented → Handle high rates and provide high-precision momentum measurements.
 - Sophisticated reconstruction algorithm containing:
 - 1. hit construction,

3.

- 2. time clustering,
- 3. tracking via pattern recognition,
- 4. refinement via Kalman fitting,
- background rejection via Machine Learning



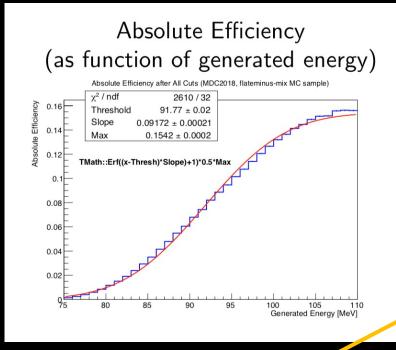
M. Devilbiss, UMich FERMILAB-SLIDES-20-100-V

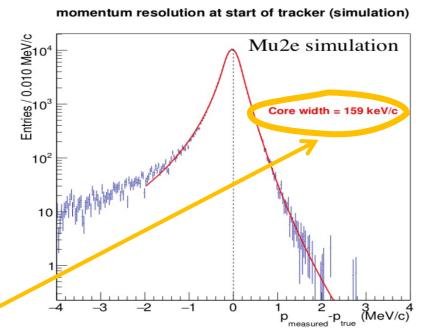


Acceptance & Response

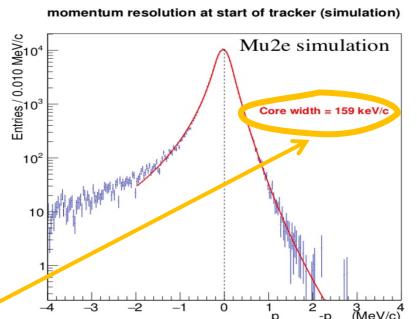
Paper documenting our Machine Learning analysis. https://arxiv.org/abs/2106.08891

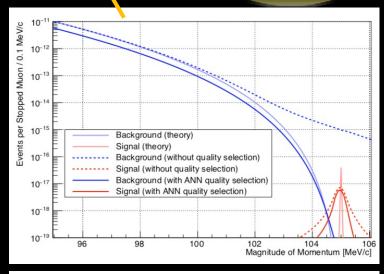
- Experiment is "blind" to anything with energy < 75 MeV.
- Tracking resolution improved by use of ANN.

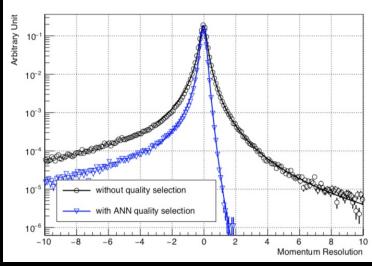




Well within our requirement!









Removing Backgrounds



Beam delivery and detector systems optimized for high intensity, pure muon beam – must be "background free":

- Intrinsic:
 - Scale with number of stopped muons.

Late arriving:

 Scale with number of late protons/ extinction performance

Туре	Source	Mitigation	Yield (over lifetime of experiment)
Intrinsic	Decay in Orbit (DIO)	Tracker Deign/ Resolution	0.144 \pm 0.028 (stat) \pm 0.11 (sys)
Late Arriving	Pion Capture	Beam Structure /Extinction	0.021 \pm 0.001 (stat) \pm 0.002 (sys)
	Pion Decay in Flight	-	$0.001 \pm < 0.001$
Other	Anti-proton	Thin Absorber Windows	0.04 \pm 0.022 (stat) \pm 0.020 (sys)
	Cosmic Rays	Active Veto System	0.209 \pm 0.0022 (stat) \pm 0.055 (sys)



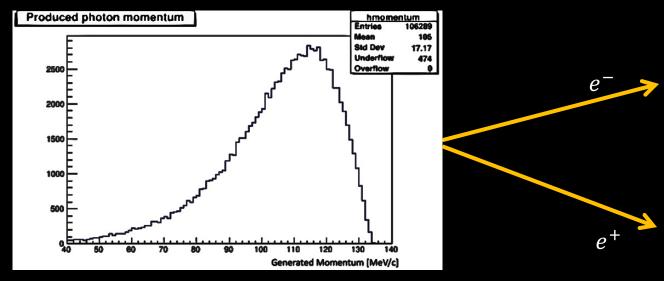
Pion Backgrounds J.A. Bistirilich, K.M. Crowe et al., Phys Rev C5, 1867 (1972)



shape templates

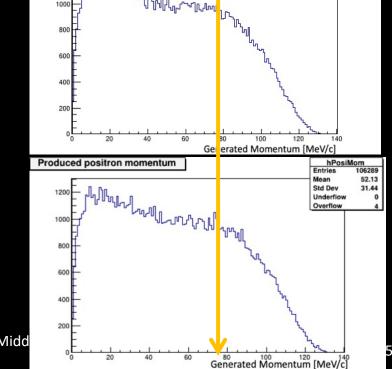
- Radiative Pion Capture occur when a pion is captured by a nucleus at our Stopping Target. The resulting photon produces an
 outgoing electron and positron pair. Pair production can be internal or external of atom.
- For simulation studies Mu2e uses the work of Bistirilich et al to parameterize the photon energy. The closest material analyzed
 in data is Magnesium. The work of Kroll-Wada-Joseph is used to calculate outgoing LorentzVectors of positron and electron:

The annular tracker means we are "blind" to large fraction of e-/e+ from RPC – even without time cuts!



Photon can be virtual (internal) or not (external)





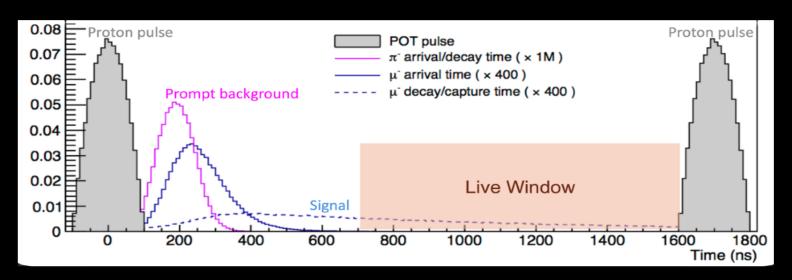
Mu2e: Searching for μ N→e N at Fermilab - Sophie Midd smidd@caltech.edu

Pion Backgrounds



Most importantly: a delayed "livegate" is enforced:

- Pions have a free lifetime of 26ns. They decay to produce muons (and neutrinos). Muons can further decay and produce backgrounds.
- Eliminate prompt backgrounds using a primary beam of short proton pulse. Use a delayed measurement window (\sim 700 ns after proton pulse at target). Before this point we ignore any tracks we see in our detector systems.



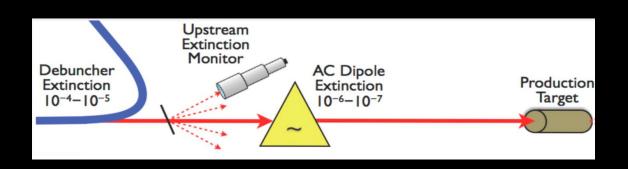
Out-of-time pions could fall inside "livegate" but are eliminated by excellent extinction in our proton beam.

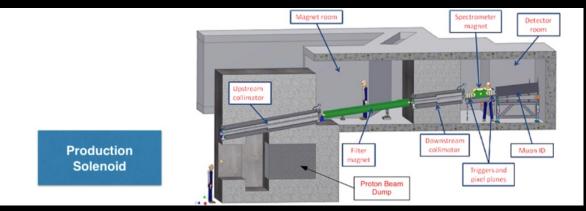


Removing out-of-time protons



- Must have out-of-time: in-time proton ratio must be kept $< 10^{-10}$ to remove potential backgrounds.
- 2 phase process:
 - Fast "kicker" which transfers the proton beam from the Recycler to the Delivery Ring preserves extinction. Extinction of 10^{-5} is expected as the proton beam is extracted and delivered.
 - The beam line from the Delivery Ring to the production target has a set of AC oscillating dipoles that sweep out-of-time protons into a system of collimators. This should achieve an additional extinction of 10⁻⁷ or better.
- Extinction measured using a detector system: Si-pixel + sampling EMC.







Removing Backgrounds



Beam delivery and detector systems optimized for high intensity, pure muon beam – must be "background free":

Туре	Source	Mitigation	Yield (over lifetime of experiment)
Intrinsic	Decay in Orbit (DIO)	Tracker Design/ Resolution	0.144 \pm 0.028 (stat) \pm 0.11 (sys)
Late Arriving	Pion Capture	Beam Structure /Extinction	0.021 \pm 0.001 (stat) \pm 0.002 (sys)
	Pion Decay in Flight	-	$0.001 \pm < 0.001$
Other	Anti-proton	Thin Absorber Windows	0.04 \pm 0.022 (stat) \pm 0.020 (sys)
	Cosmic Rays	Active Veto System	0.209 \pm 0.0022 (stat) \pm 0.055 (sys)

Scales with livetime!



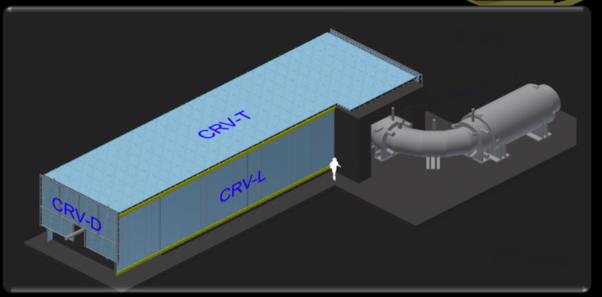
Cosmic Ray Backgrounds

Each day, ~ 1 conversion-like electron is produced by cosmic rays

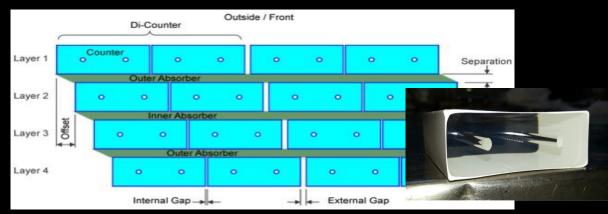


Cosmic Ray Veto will prevent cosmic muons faking a signal:

- Cosmic-ray muons can initiate 105 MeV particles that appear to emanate from the stopping target:
 - Electrons and positrons through secondary and delta-ray production in the material within the solenoids,
 - Electrons from muon decay-in-flight,
 - Muons themselves can be misidentified as electrons.
- Remove using active veto (CRV) + overburden and shielding concrete surrounding the Detector Solenoid.



Each panel is composed of 5 x 2 x 450 cm³ scintillator bars:





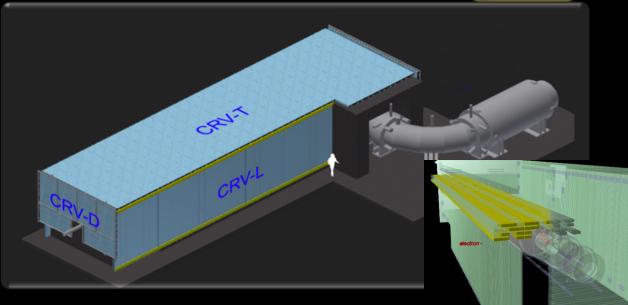
The Cosmic Ray Veto

Each day, ~1 conversion-like electron is produced by cosmic rays

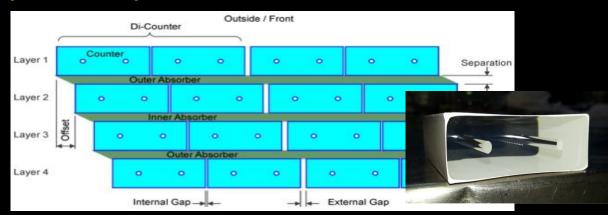


Cosmic Ray Veto will prevent cosmic muons faking a signal:

- 4 layers of extruded polystyrene scintillator counter.
- Surrounds the top and sides of DS and the downstream end of the Transport Solenoid.
- Remove Cosmic-ray candidates:
 - A track stub consisting of at least three adjacent hit strips in different layers within a 5 ns time window signals the presence of a cosmic-ray muon.
 - Signal candidates within 125 ns of such a track stub are assumed to be produced by a cosmic-ray muon and is vetoed.
 - 99.99% efficiency requirement!*
- CRV in intense radiation environment:
 - Neutrons produced at the production target, stopping target, and muon beam stop;
 - Gammas produced largely from neutron capture.
- These can make hits in CRV, faking cosmic ray muons, increasing the dead-time.



Each panel is composed of 5 x 2 x 450 cm³ scintillator bars:



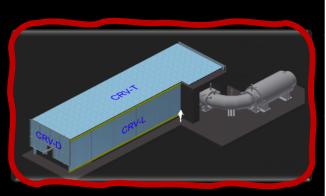


*Math documented in our TDR [1] 26 October 2021 80

Removing Backgrounds



Beam delivery and detector systems optimized for high intensity, pure muon beam – must be "background free":

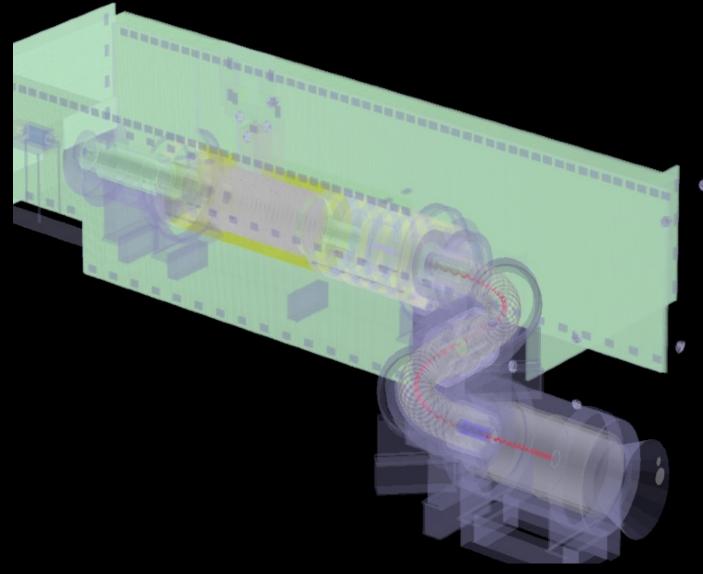


Active veto system surrounds detector region

Туре	Source	Mitigation	Yield (over lifetime of experiment)
Intrinsic	Decay in Orbit (DIO)	Tracker Deign/ Resolution	0.144 \pm 0.028 (stat) \pm 0.11 (sys)
Late Arriving	Pion Capture	Beam Structure /Extinction	0.021 \pm 0.001 (stat) \pm 0.002 (sys)
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	Cosmic Rays	Active Veto System	0.209 \pm 0.0022 (stat) \pm 0.055 (sys)







Building our Detectors

How are our detectors constructed? Where are they constructed? What is the current status of each system? When will Mu2e be ready?



The Tracker: Progress

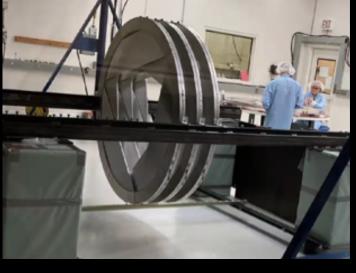
2020: Production at University of Minnesota, testing at Duke Uni., → Over 60% of panels fabricated, testing on going



→ 7/36 planes so far assembled on site











Vacuum tests at FNAL, here for single panel.

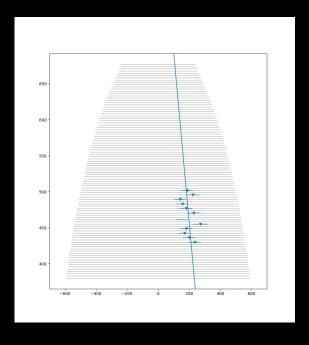




The Tracker: Progress

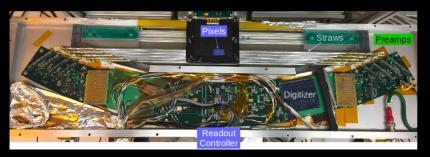
2020: Vertical Slice Test begins at FNAL



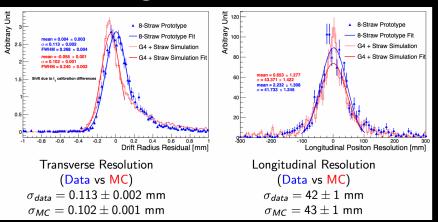


- Use Cosmic Rays.
- Use information gained to update MC.
- Measured performance and resolutions.
- First test with real data.





Measured gain, crosstalk, resolution...



8 channel prototype

Read about the prototype: → Good agreement between MC/Data https://arxiv.org/abs/1710.03799

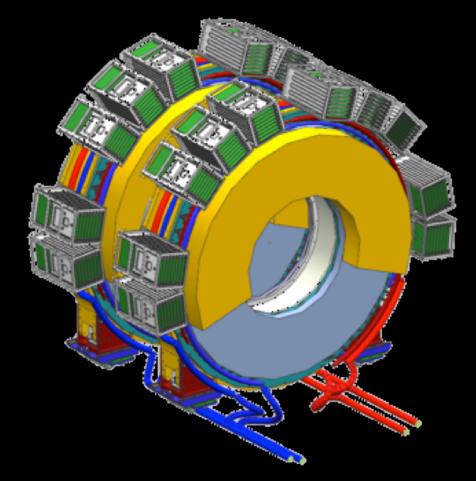
→ Resolution can be achieved



The Calorimeter: Purpose



- The calorimeter is vital for providing:
 - Particle identification,
 - Fast online trigger filter,
 - Accurate timing information for background rejection
 - Seed for track reconstruction.
- The Mu2e Calorimeter must:
 - Have a large acceptance;
 - Provide time resolution < 0.5 ns;
 - Energy resolution < 10%;
 - Position resolution of 1 cm.;
 - Function in region with radiation exposure up to 20Gy/crystal/year and with neutron flux 10¹¹ /cm².
- Annular shape and 2 disks the separation of disks is ½ the pitch of the mu2e signal helix, means that we cannot miss a conversion electron.
- Each disk = 674 CsI crystals





The Calorimeter

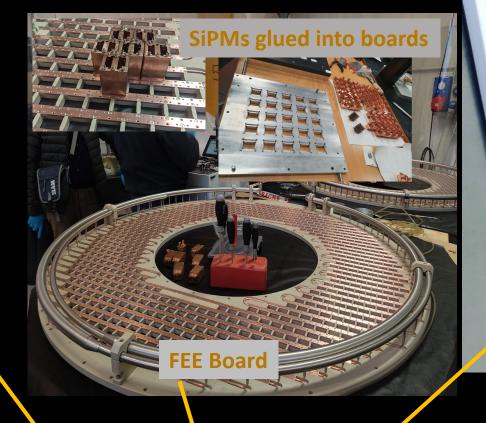
All parts ready Assembly beginning now



Inner Ring



Outer Ring

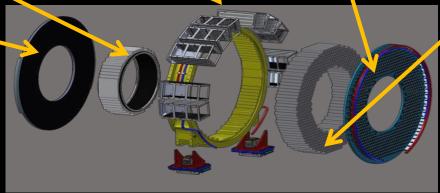


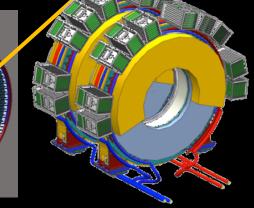
Crystals at **FNAL** in sealed cupboard after





Caltech



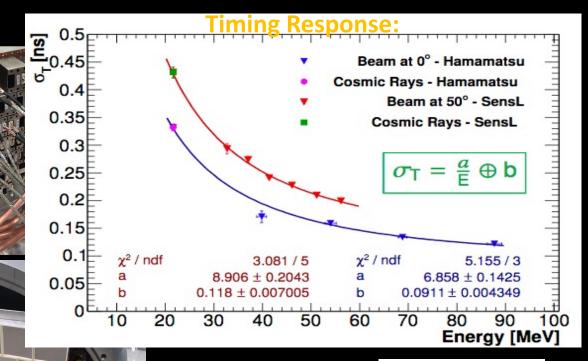


Me at FNAL finishing crystals QA (2020)26 October 202

Mu2e: Searching for μ N→e N at Fermilab - Sophie Middleton - smidd@caltech.edu

The Calorimeter: Progress

- R&D and Prototyping successfully completed.
- 51 crystals + 102 SiPM + 102 FEE boards

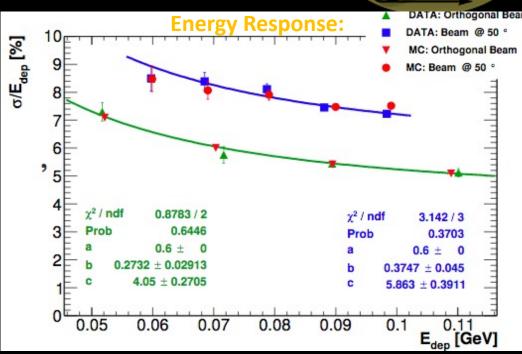


 $\frac{\text{Typical time resolution}}{\text{E}_{\text{beam}} \ @ \ 100 \ \text{MeV}}$ $\frac{\sigma_{\text{T1}} \sim 130 \ \text{ps}}{\sigma_{\text{T1}} \sim 130 \ \text{ps}}$

Read about the test beam results: https://www.osti.gov/pages/biblio/1523418

2018:





- Test beam with e^- with E = 60-120 MeV.
- Good agreement between MC/Data!
- Meets energy and timing performance requirements!

The Cosmic Ray Veto System









2020: Vertical Slice Test at UVA



- 56/83 modules fabricated at UVA
- Electronics production underway
- Front-End-Boards produced KSU
- Vertical slice test underway



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The Stopping Target Monitor (STM)



- Need an accurate measure of total number of stopped muons in the target (within 10%).
- Placed far downstream of Detector Solenoids (~34 m from target).
- STM uses HPGe and LaBr₃ detectors to measure X/gamma-rays produced by stopped muons in Al target:
 - Prompt X-ray emitted from muonic atoms at 347keV;
 - Delayed gamma ray at 844keV;
 - 3. Semi-prompt gamma ray at 1.809MeV.

HPGe and LaBr detectors procured

$$R_{\mu e} = \frac{\Gamma(\mu^- + A(Z, N) \to e^- + A(Z, N))}{\Gamma(all - captures)} < 7 \times 10^{-13} (90\% C.L)$$





Test beam at ELBE this year.... results being analyzed

2019: Test stands setup to begin DAQ development



Trigger & DAQ

Must combine information from ~500 detector data sources and apply filters to reduce the average data volume by a factor of at least 100 before it can be transferred to offline storage.

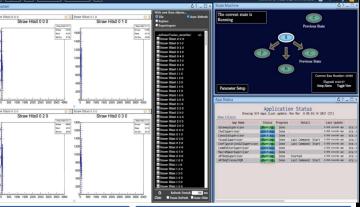
TDAQ must:

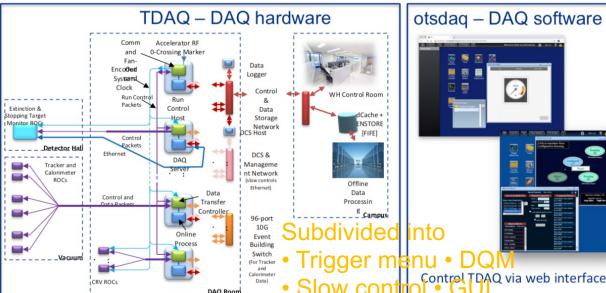
- provide efficiency > 90% on the signal;
- keep the trigger rate < a few kHz (\approx 7 PB/year);
- achieve a processing < 5 ms/event.

Event Display (REve):



DQM being developed:





Infrastructure detailed in: https://arxiv.org/abs/2010.16208



Trigger & DAQ: Progress

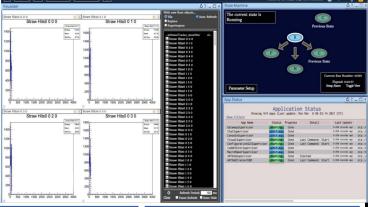
Event Display (REve):

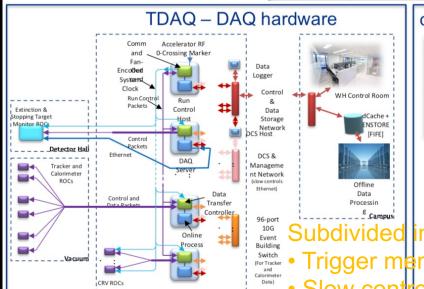


Lots of progress over last few years:

- In 2019: Test stand set up at FNAL
 - A joint platform (OTSDAQ) has been set up to allow compatibility between trackers, calorimeter and STM interfaces.
- In 2020: DQM development, event display developed using cutting edge EVE-7, event builder prototyped.
- In 2021: Horizontal slice test this summer.

DQM being developed:







Infrastructure detailed in: https://arxiv.org/abs/2010.16208



Mu2e: Timeline





End of 2024 – onwards

- → Physics data taking:
- Run 1: Until 2026 (O(10⁻¹⁶)).
- Run2: After LBNF Shutdown (O(10⁻¹⁷)).



Run 1:

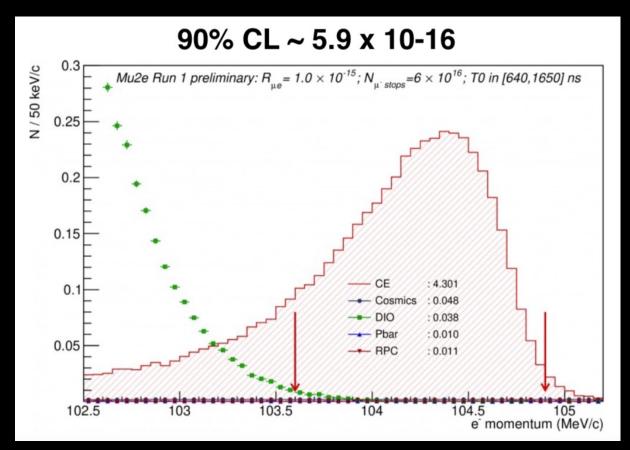
Beam on target in late by 2024

Run1: 2025-2026:

- 10³ improvement over SINDRUM-II 90% CL limit
- PIP-II/LBNF shutdown scheduled for end of 2026

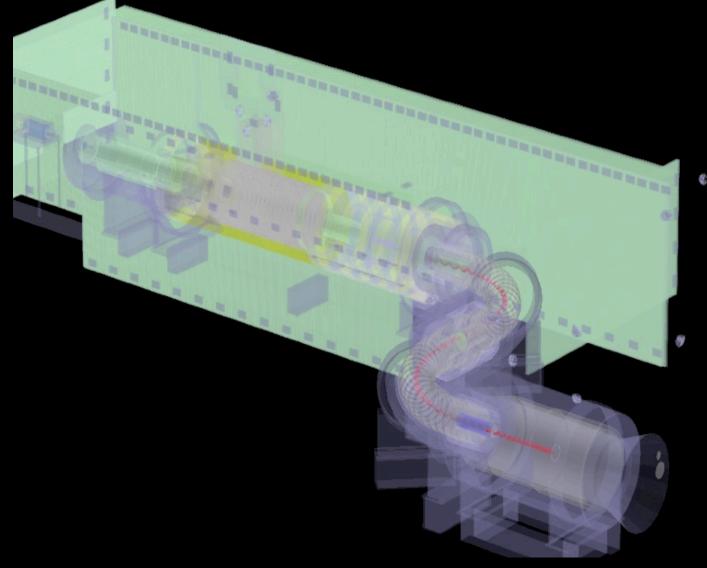
Run2: Data-taking resumes early 2029

• The goal is a x10⁴ improvement over SINDRUM-II: (90% CL)









Extended Mu2e into the PIP-II Era: Mu2e-II

What happens if we see at Mu2e signal? What happens if we don't?



Mu2e-II

Mu2e-II aims to improve the sensitivity ($R_{\mu e}$) to the neutrinoless conversion of a muon-to an-electron in the field of a nucleus by a further order of magnitude than Mu2e i.e. SES $\sim \vartheta(10^{-18})$



- There are 2 possible outcomes from Mu2e:
 - 1. Conversion not observed motivates pushing to higher mass scales .
 - **Conversion observed** motivates more precise measurements with different targets.
- Either way Mu2e-II is well motivated!

Mu2e-II would:

- Be based at Fermilab. Will utilize the (nominal) 100kW beam from Proton Improvement Plan II (PIP-II).
- Start a few years after the end of Mu2e run with an expected 3+1 years of physics running.
- Salvage and refurbish as much of Mu2e infrastructure as possible.
- Upgrade Mu2e components where required to handle higher beam intensity.
- Mu2e-II has a support from muon physics community and Fermilab's PAC
- If approved, Mu2e-II expects to start data taking at the end of the decade
- Large effort to write White Paper for Snowmass 2022!

See here for overview: https://arxiv.org/pdf/1802.02599



The PIP-II Project

- The project received CD-1 approval from the U.S. Department of Energy in July 2018.
- PIP-II will power both DUNE and other experiments like Mu2e-II.
- PIP-II is planned to deliver beam in the next decade.
- Groundbreaking ceremony took place in 2019.

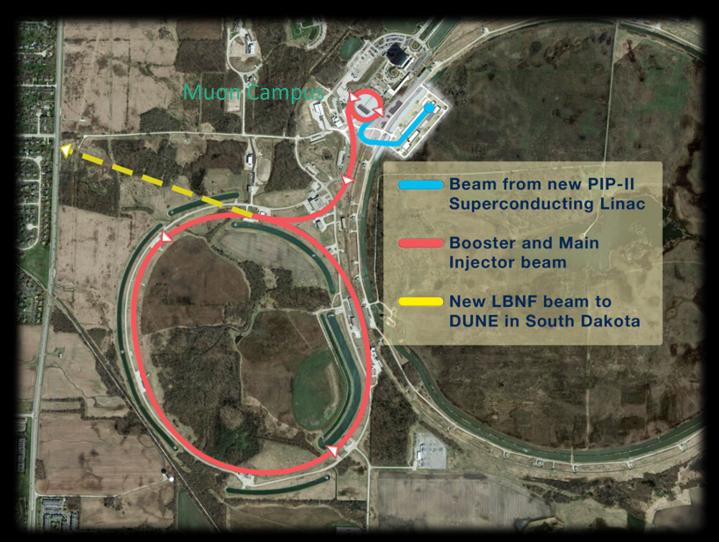






PIP-II





- PIP-II designed to deliver 800 MeV H- beam to the Booster.
- Capable of running in CW mode with 2 mA average current at 1.6 MW.
- Mu2e-II will get a beam at upstream end of transfer line to Booster:
 - Need to build a beamline to deliver beam to M4 enclosure



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Production Target



LDRD Project on-going to investigate production target choice

	Tungsten/WC	Lower-density bent (Carbon)
Rotated	Requires a large hardware in HRS	Too large to fit HRS
Fixed granular	DPA is too high	DPA is high; lower pion production
Conveyor	Thermal analysis is ongoing	Lower pion production; thermal analysis is ongoing

Front runner is Conveyor design. But made out of W or C?

Prioritizing designs

Constraint: compatibility with the current HRS design (inner bore=20 (25) cm)

Rotator

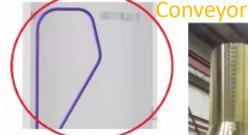


Pros: radiation damage can be Pros: small space required distributed over many rods Cons: its hardware would require a significant space inside the bore (complicates cooling and muon flow)

Granular



Cons: peak DPA (MARS15) >300/yr; gas cooling cannot be performed efficiently



Pros: small space required; He gas could be used for both cooling and moving elements inside conveyor; radiation damage can be distributed; Cons: technical complexity (prototyping needed)





Mu2e: Searching for μ N→e N at Fermilab - Sophie Middleton smidd@caltech.edu

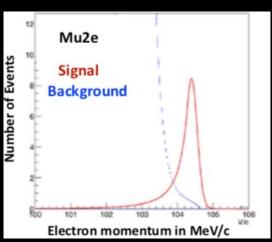
(prototyping needed) technical complexity

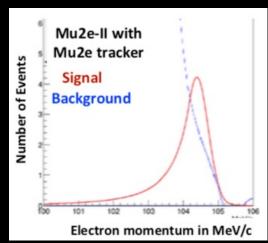
26 October 2021

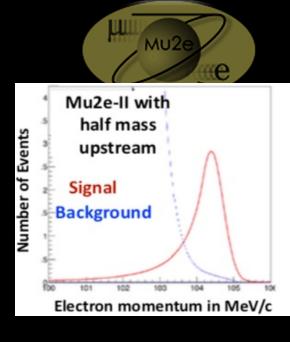
Tracker Requirements

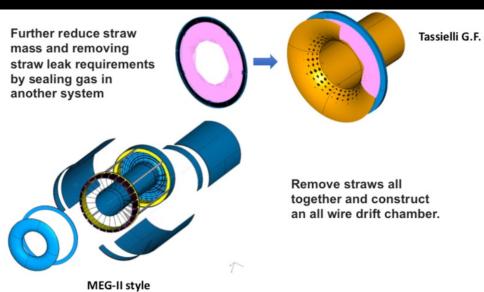
DIO background would increase x10 in Mu2e-II.

Must improve momentum resolution to suppress DIO.









To meet Mu2e-II momentum resolution/background separation goals: Reduce total Tracker Mass:

- Thinner straws (8 μ m)
- Remove the 200 angstrom gold layer from inside straw

Change detector design:

- Use an ultra light gas vessel to ease straw leakage requirements
- Use different gas
- Consider all wires construction and remove the straws
- Or wires separated by mylar walls

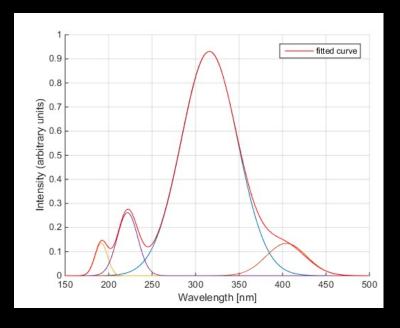
Increased hit occupancy and timing window:

4x increase in PBI is estimated to reduce reconstruction efficiency by 30%.

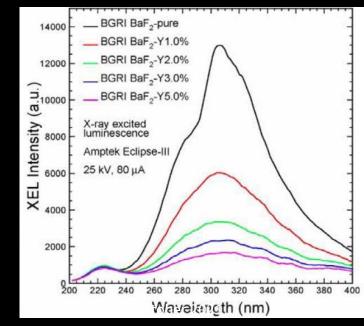
Barium Fluoride Calorimeter Crystals

- Radiation doses and rates at Mu2e-II are high for CsI:
 - Up to 900 krad and 1E13 n [1MeVeq/cm2]
- BaF₂ is an excellent candidate for a fast, high rate, radiation-hard crystal for the Mu2e-II calorimeter:
 - BaF₂ can survive up 100 Mrad
- Must have way of utilizing 220 nm fast component without interference from the larger 320 nm slow component.
- Slow suppression achieved by:
 - 1. Rare Earth Doping (Y, La,Ce).
 - 2. Develop photo-detectors sensitive to UV only:
 - SiPM with an external filter
 - UV-sensitive photocathodes
 - Solar-blind MCP SiPM "sees" only fast component.

R&D Collaboration between Caltech, JPL & FBK on-going



Actually two fast components (t = 0.6 ns) at 195 and 220 nm and two slow components (t = 630 ns) at 320 and 400 nm.

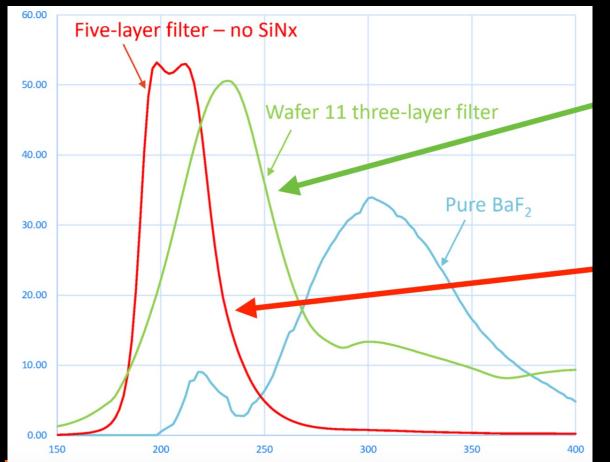


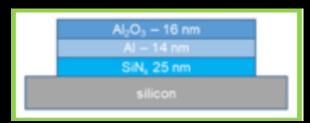


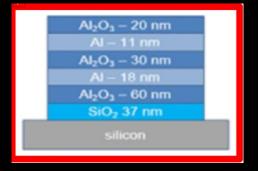
SiPM R&D



- Caltech-JPL-FBK consortium working on delivering developing a special locating for SiPMs
- Sandwich of Al, SiN_2 and Al_2O_3 layers deposited on the active material







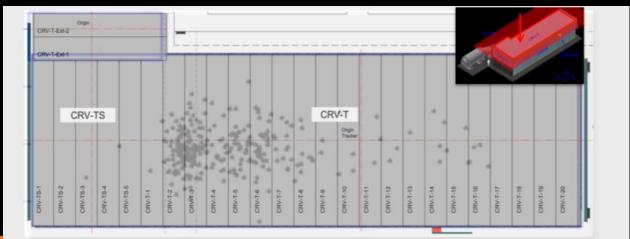


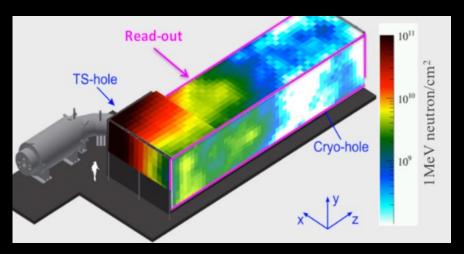
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The Cosmic Ray Veto System (CRV)



- Expected live-time and therefore Cosmic Ray backgrounds will be 3 x higher for Mu2e-II
 - Need to enhance the CRV performance in critical regions
- Light Yield degradation impacts CRV performance
 - Must replace CRV
 - Higher noise rates (x2-3) these impose challenges:
 - Higher DAQ rates
 - Radiation damage
 - Induced dead-time
 - → Enhanced shielding, fine-granular layers, other technologies



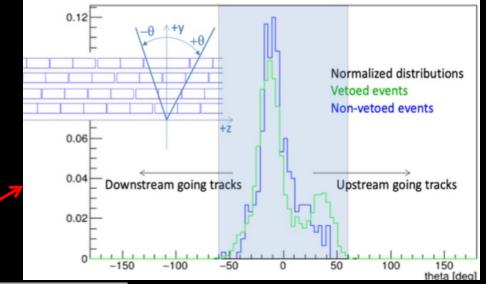


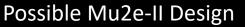


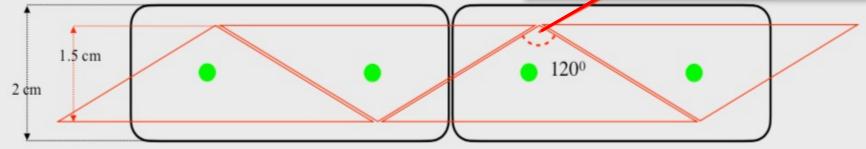
The Cosmic Ray Veto System



- Gaps between modules and counters and modules impact the CRV performance:
 - Reduce Gaps
 - Change geometry
 - Extra Layers
- Triangular Bars:
 - Improved efficiency due to reduced gaps
 - Lower dead time
 - Lower rate per channel
 - Simple design









Trigger & Data Acquisition (TDAQ)



Increased data rate, more background and more detector channels:

- 10x data rate
- X3 event size
- 3000:1 rejection is needed to arrive at 14PB/year

Considerations:

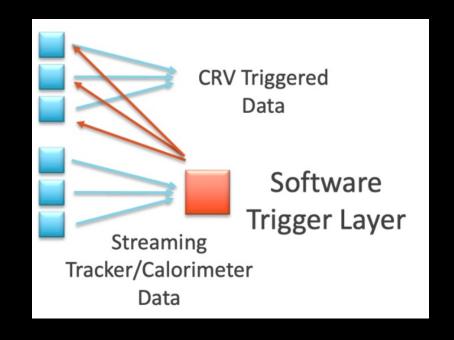
- Reduced off-spill time to readout large front-end buffers
- Streaming .v. triggered data taking
- Radiation tolerances requirements

No large buffers for the CRV:

- Large CRV buffers + software trigger
- Small CRV buffers + hardware trigger

Solutions:

- 2-level TDAQ based on FPGA pre-processing and trigger primitives
- 2-level TDAQ system based on FPGA pre-filtering
- TDAQ based on GPU co-processor
- Trigger-less TDAQ based on software trigger.





Follow us





Watch the experiment evolve with frequent videos and images:



https://twitter.com/Mu2eExperiment



https://www.instagram.com/mu2eexperiment/



Summary



- Muon CLFV channels offer deep indirect probes into BSM.
- Mu2e is at the forefront of active global CLFV program. Discovery potential over a wide range of well motivated BSM models.
- Muon-to-electron sector complements tau and Higgs collider searches such as: $\tau \to e \gamma$ or $\mu \gamma$ and $H \to e \tau$, $\mu \tau$, or μe .
 - It is important to eliminate Standard Model backgrounds so the experiment is designed to be "background free":
 - Super conducting solenoids to collect and efficiently transport low momentum muons;
 - Pulsed beam removes backgrounds from pions;
 - Low mass, annular tracker has high resolution to avoid DIO backgrounds;
 - Cosmic Ray Veto surrounds detectors to remove "fake signals" from Cosmic muons.
- Looking further ahead the Mu2e-II experiment will help elucidate any signal and push to higher mass scales (of no signal).
- Large effort on-going to design the Mu2e-II experiment with consideration of requirements introduced due to higher beam intensity.

Thank You for listening!



Useful Resources



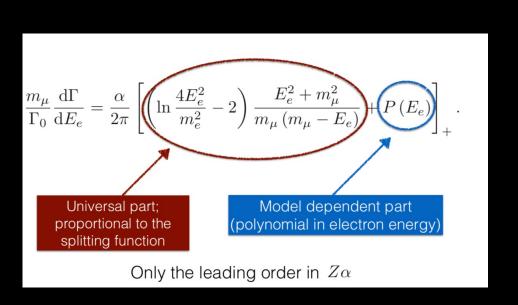
- 1. S. T. Petcov, Sov. J. Nucl. Phys. **25**, 340 (1977); Yad. Fiz. **25**, 1336 (1977) [erratum].
- 2. S. M. Bilenky, S. T. Petcov, and B. Pontecorvo, Phys. Lett. B **67**, 309 (1977).
- 3. W. J. Marciano and A. I. Sanda, Phys. Lett. B **67**, 303 (1977).
- 4. B. W. Lee, S. Pakvasa, R. E. Shrock, and H. Sugawara, Phys. Rev. Lett. **38**, 937 (1977); **38**, 1230 (1977) [erratum].
- 5. J. Adam *et al.* (EG Collaboration), Phys. Rev. Lett. **110**, 20 (2013).
- 6. W. Bertl *et al.* (SINDRUM-II Collaboration), Eur. Phys. J. **C47**, 337 (2006).
- 7. U. Bellgardt *et al.*, (SINDRUM Collaboration), Nucl. Phys. **B299**, 1 (1988).
- 8. A.M. Baldini *et al.*, "MEG Upgrade Proposal", arXiv:1301.7225v2 [physics.ins- det].
- 9. Y. Kuno et al., "COMET Proposal" (2007) see also https://arxiv.org/abs/1812.09018 for Phase I TDR
- 10. Mu2e TDR, arXiv:1501.05241
- 11. Nuclear Physics B Proceedings Supplements Volumes 248–250, March–May 2014, Pages 35-4
- 12. A. Czarnecki et al., "Muon decay in orbit: Spectrum of high-energy electrons," Phys. Rev. D 84 (Jul, 2011).
- 13. Sindrum-II "Improved limit of Branching Fraction of mu- \rightarrow e+ in Titanium", Phys Lett B 422 (1998) 334-338 (1998)

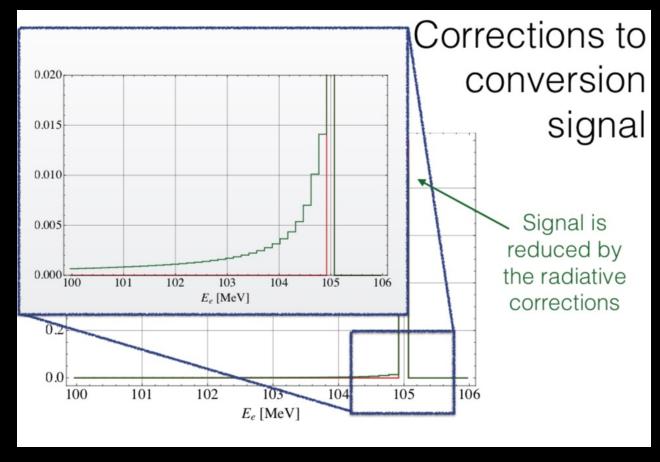


Radiative Corrections for CE Signal



Our signal also requires radiative corrections:







g-2 Result: Implications for Mu2e



P. Paradisi / Nuclear Physics B (Proc. Suppl.) 248-250 (2014)

• Dipole transitions $\mu \to e \gamma$ in the leptonic sector are accounted for by means of the effective Lagrangian :

$$\mathcal{L} = e \frac{m_{\ell}}{2} \left(\bar{\ell}_R \sigma_{\mu\nu} A_{\ell\ell'} \ell'_L + \bar{\ell}'_L \sigma_{\mu\nu} A^{\star}_{\ell\ell'} \ell_R \right) F^{\mu\nu},$$

$$\frac{\mathrm{BR}(\ell \to \ell' \gamma)}{\mathrm{BR}(\ell \to \ell' \nu_{\ell} \bar{\nu}_{\ell'})} = \frac{48\pi^3 \alpha}{G_F^2} \left(|A_{\ell\ell'}|^2 + |A_{\ell'\ell}|^2 \right).$$

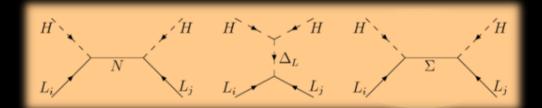
- The underlying $\mu \to e \gamma$ transition can also generate lepton flavor conserving processes like the anomalous magnetic moments (Δa_{μ}) as well as leptonic electric dipole moments (EDMs, d_{μ}).
- In terms of the effective Lagrangian can write as:

$$\Delta a_{\ell} = 2m_{\ell}^2 \operatorname{Re}(A_{\ell\ell}), \qquad \frac{d_{\ell}}{e} = m_{\ell} \operatorname{Im}(A_{\ell\ell}).$$

On general grounds, one would expect that, in concrete NP scenarios, (Δa_{μ}) , d_{μ} and BR($\mu \to e \gamma$), are correlated. In practice, their correlations depend on the unknown flavor and CP structure of the NP couplings and thus we cannot draw any firm conclusion that we would necessary see CLFV in the next generation, but this is of course a very promising result for muon physics!



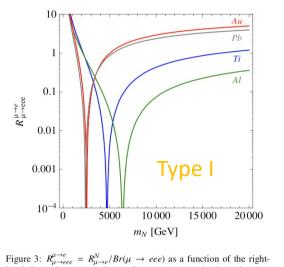
Example: See Saw Mechanisms

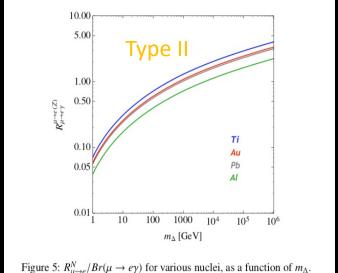


- See Saw Models can induce rates which are not suppressed by smallness of these masses.
- There are 3 ways of inducing $\Delta L = 2$ Majorana neutrino masses from the tree level exchange of a heavy particle:
 - Type I exchange of right-handed neutrinos N_i,
 - Type II exchange scalar triplet ΔL_{r}
 - Type II exchange of fermion triplets Σ_i .

Knowledge of the neutrino mass matrix is not sufficient to be able to distinguish between the 3 seesaw models \rightarrow CLFV can

help here.







SUSYSO(10)

Complementary to Muon g-2 and LHC program:

SUSY SO(10)

Consider SO(10) SUSY GUT model with very massive right-handed neutrinos. Can consider different hypothesis for the neutrino Yukawa couplings. Mu2e will be able to test all PMNS type and most CKM type SO(10).

L. Calibbi et al., JHEP 1211 (2012) 040

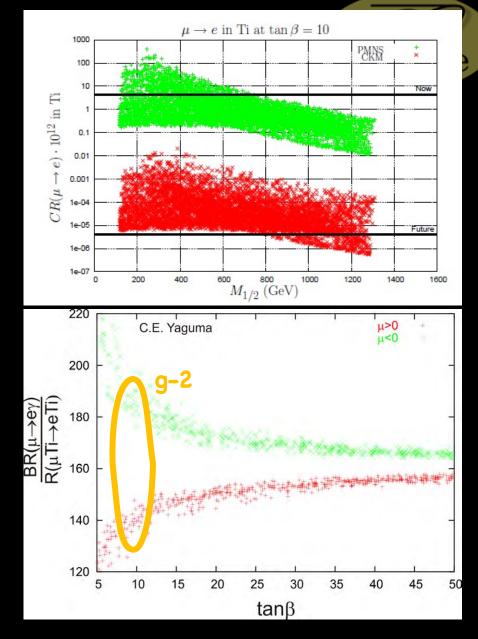
L. Calibbi, G. Signorelli arXiv:1709.00294

Muon g-2 results will also be helpful here.

To allow discrimination among different models Need:

- Observation of CLFV in more than one channel, and/or
- Evidence from LHC and/or g-2

Yaguna, hep-ph/0502014v2 Endo arxiv.org/abs/1303.4256v1 (g-2 SUSY .v. LHC constraints)





COMET: Phased Implementation



