Eigenstate Thermalization in Sachdev-Ye-Kitaev model & Schwarzian theory

Pranjal Nayak

[basesd on 1901.xxxxx]

with Julian Sonner and Manuel Vielma



Closed Quantum Systems

Quantum Mechanics is unitary!

$$|\psi(t)\rangle = \mathcal{U}(t,t_0)|\psi(t_0)\rangle$$

How come we observe thermal physics?

How come we observe black hole formation?

Plan of the talk

- √ Review of ETH
- √ Review of Sachder-Ye-Kitaer model
- √ Numerical studies of ETH in SYK model
- √ ETH in the Schwarzian sector of the SYK model
- VETH in the Conformal sector of the SYK model
- V A few thoughts on the bulk duals
- √ Summary and conclusion

Eigenstate Thermalization Hypothesis (ETH)

Classical thermalization: **ergodicity** & ← chaos

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V.S.

Quantum thermalization: Eigenstate Thermalisation Hypothesis

$$\langle m|\mathcal{O}|n\rangle = \overline{\mathcal{O}}_{\mathrm{mc}}(\overline{E})\delta_{mn} + e^{-S(\overline{E})/2}f(\overline{E},\omega)R_{mn}$$

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micro-canonical ensemble avg. energy of ensemble

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micro-canonical entropy 1 in RMT

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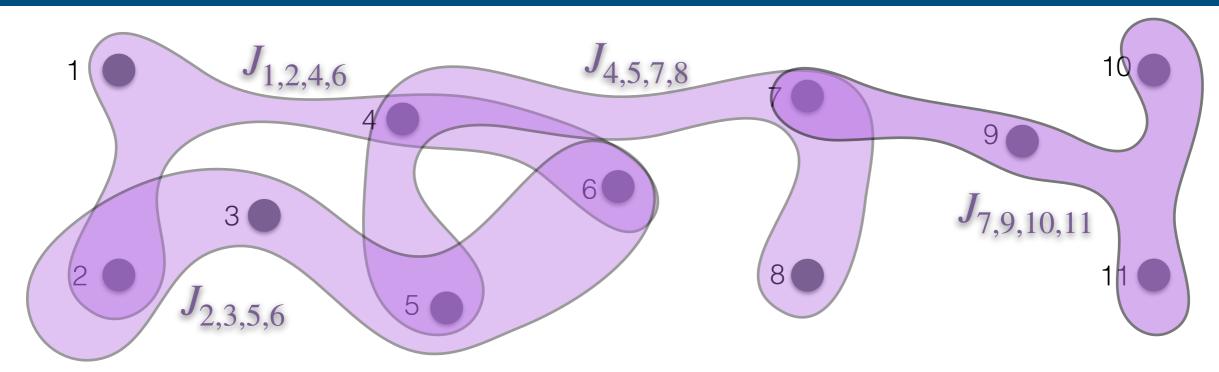
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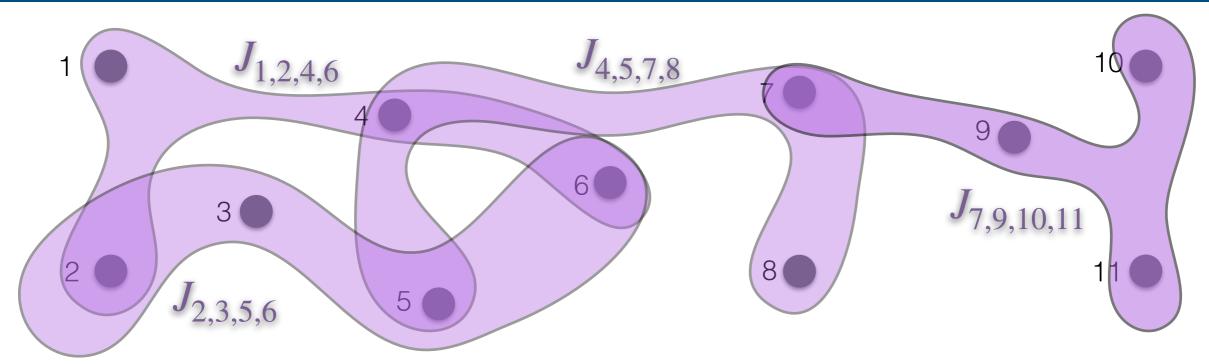
A generic excited state will then thermalise by dephasing

$$\langle \psi | \mathcal{O} | \psi \rangle = \sum_{i,j} c_i^* c_j e^{it(E_i - E_j)} \mathcal{O}_{ij} \longrightarrow \overline{O}\left(\bar{E}\right) + e^{-S}$$
 expectation value of non-extensive operator dephasing: spectral chaos on average thermal up to exponential in S

Sachder-Ye-Kilaer (SYK) model: a Review

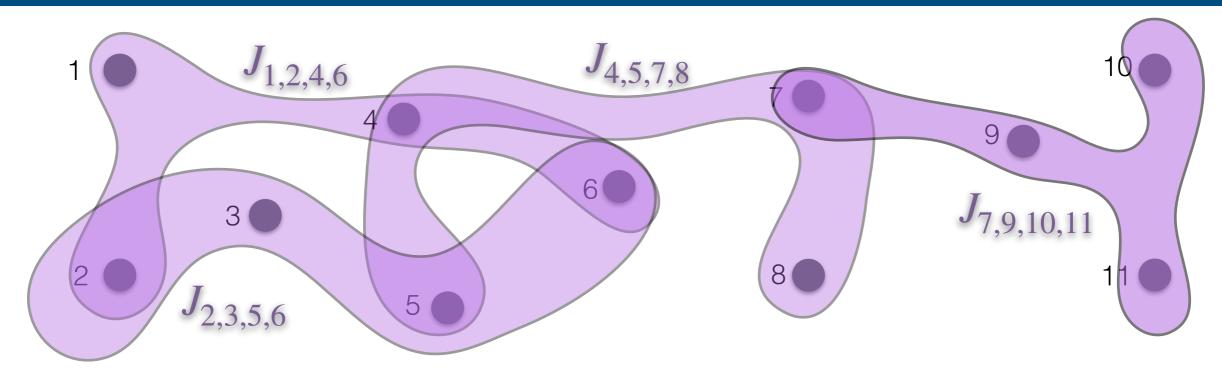


[Sachdev-Ye '93; Kitaev '15]



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- with all-to-all couplings
- > and quenched random couplings



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$$H = -\sum_{1 \le i_1 < i_2 < i_3 < i_4 \le N} J_{i_1 i_2 i_3 i_4} \psi_{i_1} \psi_{i_2} \psi_{i_3} \psi_{i_4}$$

$$1 \le i_1 < i_2 < i_3 < i_4 \le N$$

where, J is chosen from a Gaussian ensemble:

$$\langle J_{ijkl} \rangle = 0$$
 $\langle J_{ijkl}^2 \rangle = \frac{3! J^2}{N^3}$

[Sachdev '15; Parcollet, Georges '00; Kitaev '15; Polchinski, Rosenhaus '16; Jevicki, Suzuki, Yoon '16; Maldacena, Stanford '16]

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at large-N, it is often useful to use bi-local effective action

$$-\frac{S_{\text{col}}}{N} = \log \operatorname{Pf}\left[\partial_{\tau} - \Sigma\right] + \frac{J^{2}}{2q} \int d\tau d\tau' |G(\tau, \tau')|^{q} - \frac{1}{2} \int d\tau d\tau' \Sigma(\tau', \tau) G(\tau, \tau')$$

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$$J^2 G^{q-1}(\tau, au') = \Sigma(au, au')$$

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captures Melonic Physics!

Emergent Symmetry of SYK

[Kitaev '15; Polchinski, Rosenhaus '16; Jevicki, Suzuki, Yoon '16; Maldacena, Stanford '16]

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▶ The action and the saddle point equations are symmetric under reparametrizations:

$$G(\tau, \tau') \to \left(f'(\tau) f'(\tau') \right)^{2/q} G\left(f(\tau), f(\tau') \right)$$

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▶ This symmetry is explicitly broken by considering the 1/J corrections.

schwarzian in syk

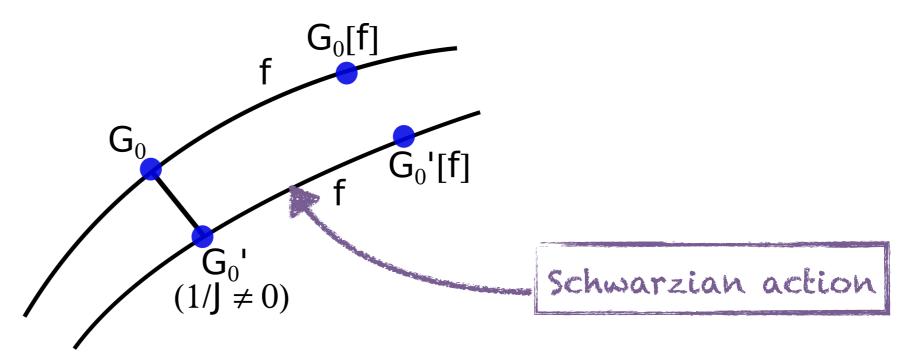
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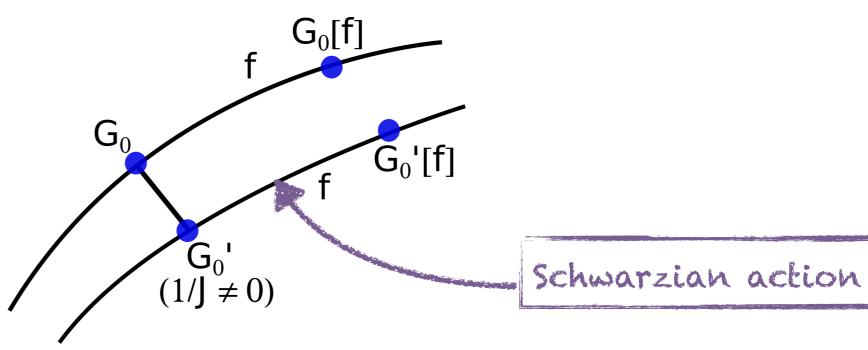
mode physics



schwarzian in syk

[Kitaev '15; Maldacena, Stanford '16]

At low-energy break conformal symmetry. Leading softmode physics



Effective action on the 'reparametrization modes'

$$\int \frac{f(\tau)}{\mathbb{SL}(2,\mathbb{R})} \exp\left[-\frac{1}{g^2} \int d\tau \left\{ f(\tau), \tau \right\} \right]$$
where, $\left\{ f(\tau), \tau \right\} = \frac{f'''(\tau)}{f'(\tau)} - \frac{3}{2} \left(\frac{f''(\tau)}{f'(\tau)}\right)^2$

$$g^2 \sim \frac{\beta J}{N}$$

[Kitaev '15; Polchinski, Rosenhaus '15; Jevicki, Suzuki, Yoon '16; Maldacena, Stanford '16; Gross, Rosenhaus '17]

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discrete tower of states

$$\mathcal{O}_n \sim \psi_i \, \partial^{2n+1} \psi_i$$

$$h_n = 2n + 1 + \epsilon_n$$

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continuum: Schwarzian

$$S = -\frac{1}{g^2} \int d\tau \left\{ f(\tau), \tau \right\}$$

$$f(\tau) \in \mathbf{Diff}(S^1)$$

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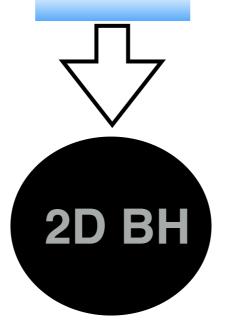
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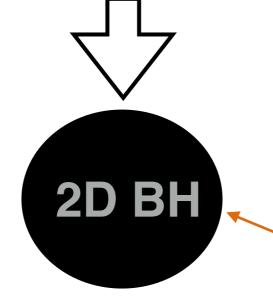
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Limit of conformal six-pt functions OPE coeffs.

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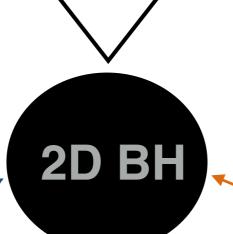
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Ward identities for Schwarzian ~ monodromy method

weak ETH



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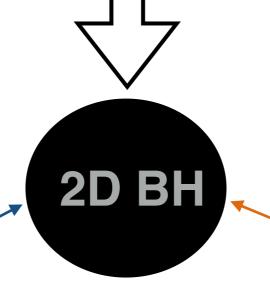
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ETH

ETH in the 57K model: A numerical study

Study by Exact Diagonalization

▶ Solve SYK in exact diagonalization [Sonner, Vielma '17]

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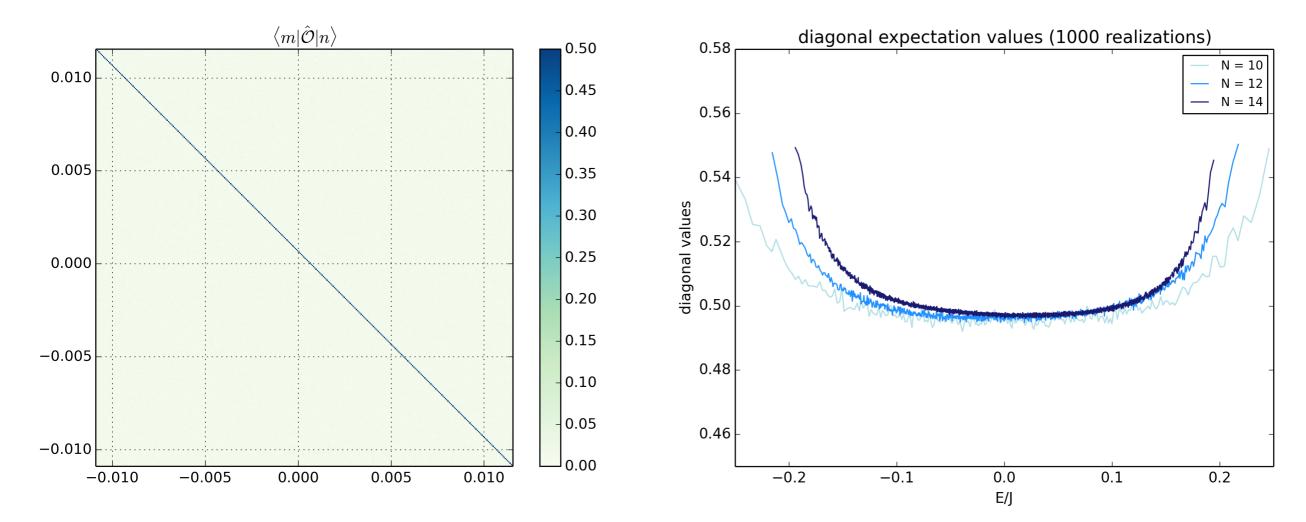
number operator, \hat{n}_k , at some site k

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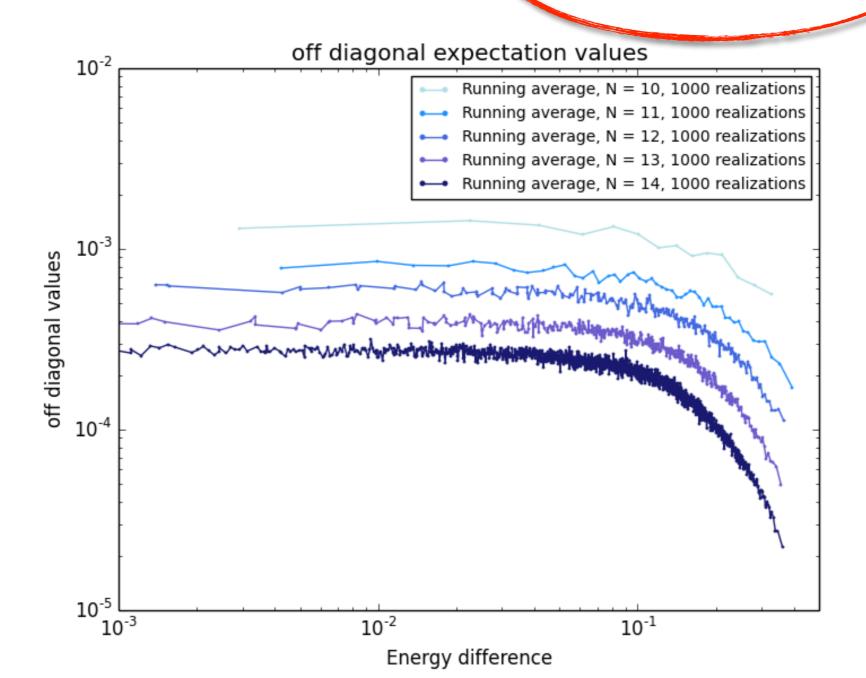
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CMT or Not?

Need to look at the off-diagonal matrix elements

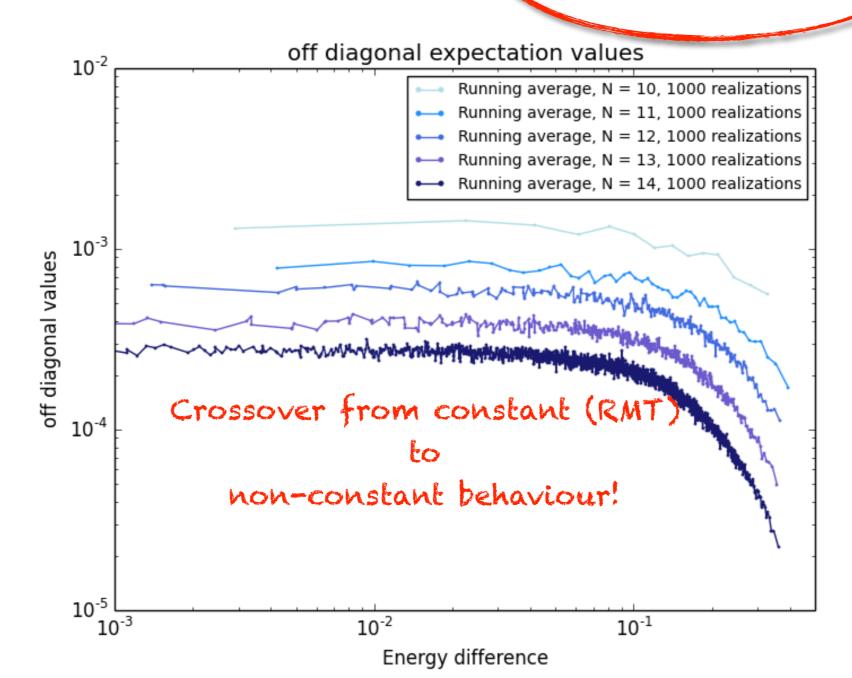
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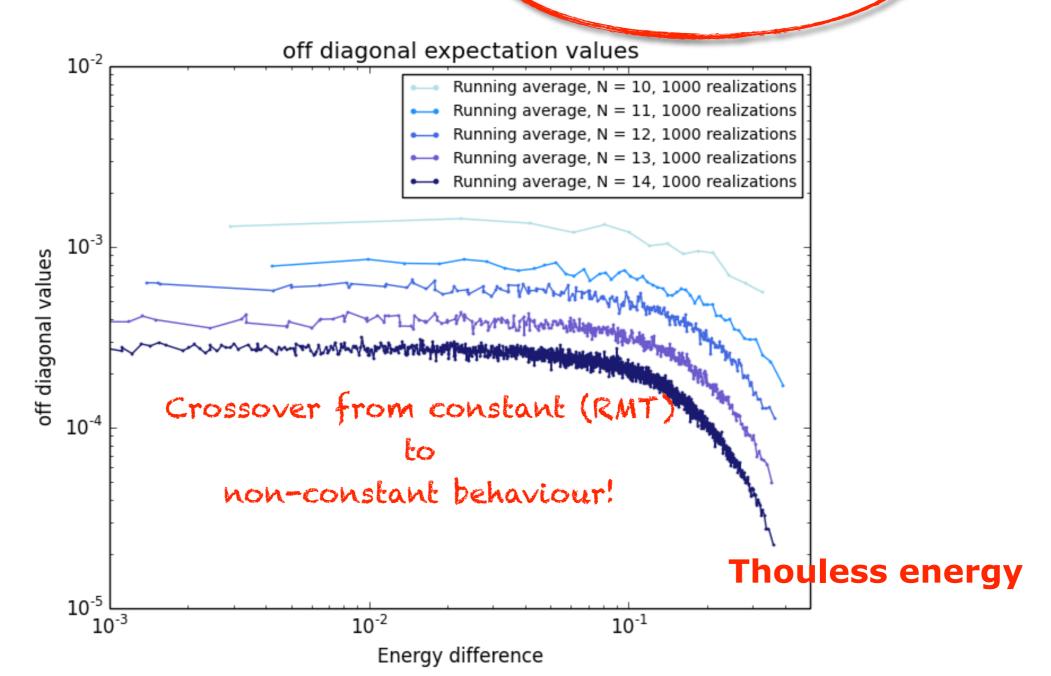
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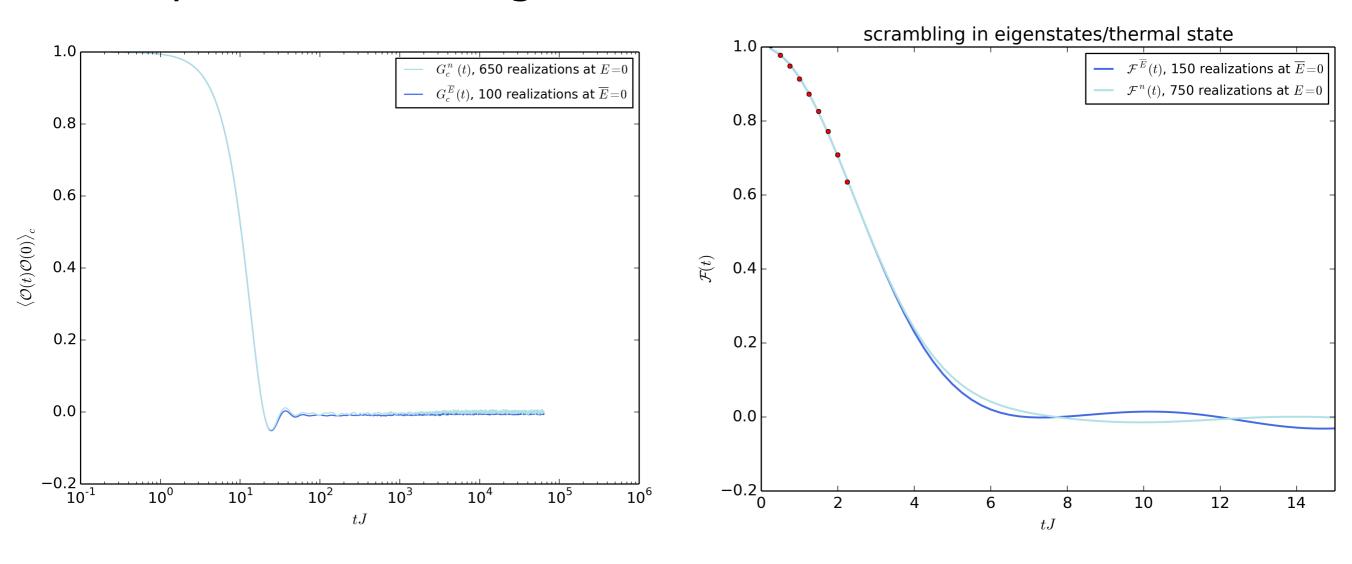
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More evidence of ETH

[Sonner, Vielma '17]

Compare OTOC in eigenstates to thermal result



Become essentially indistinguishable as system size increases

Conjecture:
$$\exists \lambda_{\rm L}^{\rm ETH} = \frac{2\pi}{\beta(E)}$$

? Can we understand the numerically observed thermalization in a finer detail? Analytically?

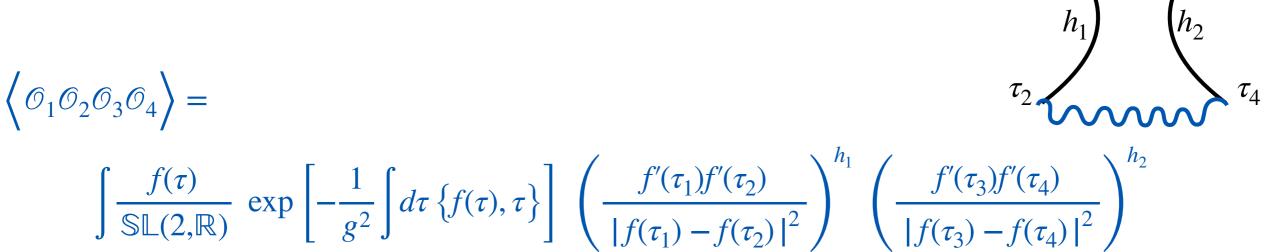
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- ? Can we understand how and why the results deviate from RMT behaviour?
- ? Can we do a dual bulk computation to understand black hole formation in 2 dimensions?

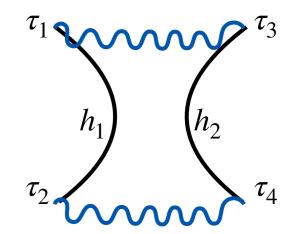
ETH in the Schwarzian Limit

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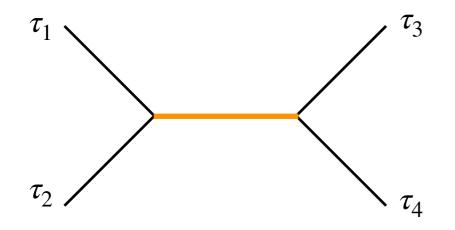
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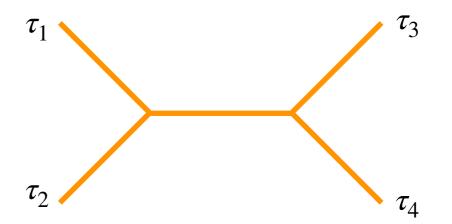


$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \mathcal{O}_4 \rangle =$$

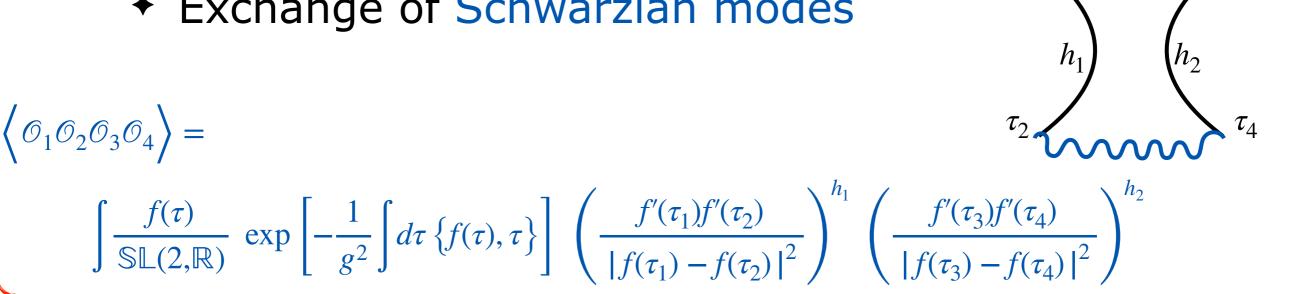
$$\int \frac{f(\tau)}{\mathbb{SL}(2,\mathbb{R})} \exp \left[-\frac{1}{g^2} \int d\tau \left\{ f(\tau), \tau \right\} \right] \left(\frac{f'(\tau_1) f'(\tau_2)}{|f(\tau_1) - f(\tau_2)|^2} \right)^{h_1} \left(\frac{f'(\tau_3) f'(\tau_4)}{|f(\tau_3) - f(\tau_4)|^2} \right)^{h_2}$$

Exchange of excited modes (conformal limit)

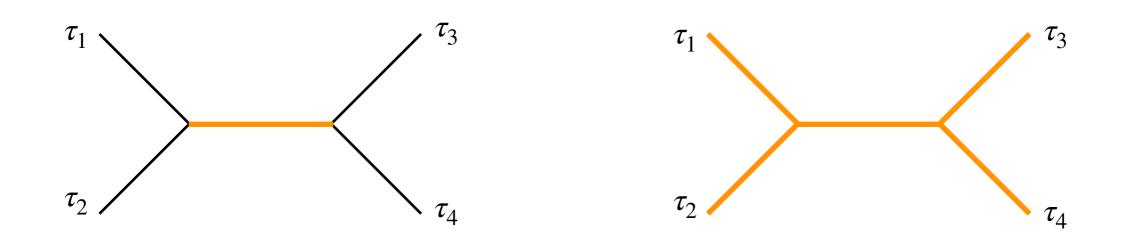




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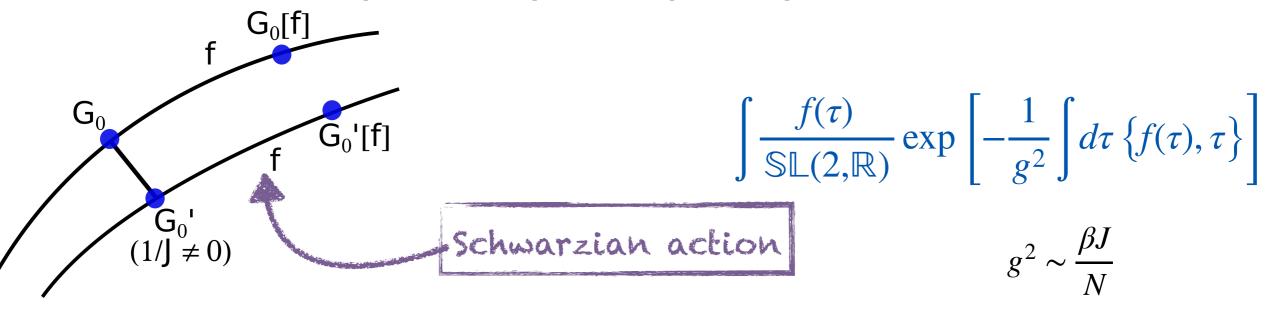


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The Schwarzian Enecry

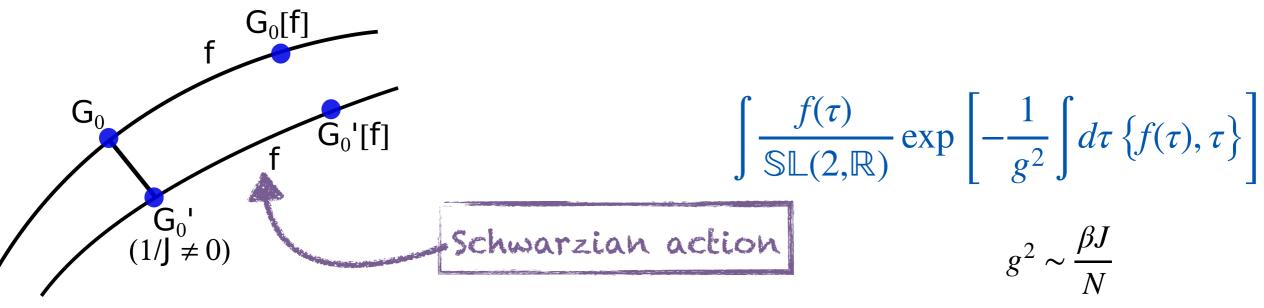
The Schwarzian Encory

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• $f(\tau)$ are the pseudo-Goldstone modes

The leading contribution to the physical observables is due to exchange of these modes

$$\left\langle \cdot \right\rangle = \int \frac{f(\tau)}{\mathbb{SL}(2,\mathbb{R})} \exp \left[-\frac{1}{g^2} \int d\tau \left\{ f(\tau), \tau \right\} \right] \left(\cdot \right)$$

Ward Identities

[PN, Sonner, Vielma '19]

Ward Identifies

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 Within Schwarzian theory, one can derive the following 'Ward' identity

$$\sum_{i=1}^{m} \left[h_i \left(\delta(x - x_{2i-1}) + \delta(x - x_{2i}) \right) + \left(sgn(x - x_{2i-1}) - sgn(x - x_{2i}) \right) \partial_{x_{2i}} \right] \prod_{j=1}^{n} \left\langle \mathcal{O}_j \left(x_{2j-1}, x_{2j} \right) \right\rangle$$

$$= -\frac{1}{g^2} \langle Sch(x) \prod_{i=1}^n \mathcal{O}_i(x_{2j-1}, x_{2j}) \rangle + \kappa_3$$

$$\mathcal{O}_j(x_{2j-1}, x_{2j}) \equiv \mathcal{O}_{h_j}(x_{2j-1}) \mathcal{O}_{h_j}(x_{2j})$$

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- We are interested in computing $\langle HHLL \rangle$ where, 'heavy' operators are defined by $h_i \sim 1/g^2$
- In the semi-classical limit, $g^2 \rightarrow 0$,

$$\{f(\tau), \tau\} = -g^2 \sum_{i} \left[\kappa_{1,i} \delta(\tau - \tau_{2i-1}) + \kappa_{1,i} \delta(\tau - \tau_{2i}) - \kappa_{2,i} \left(sgn(\tau - \tau_{2i-1}) - sgn(\tau - \tau_{2i}) \right) \right] + g^2 \kappa_3$$

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$$\mathcal{O}_j(x_{2j-1}, x_{2j}) \equiv \mathcal{O}_{h_j}(x_{2j-1}) \mathcal{O}_{h_j}(x_{2j})$$

- We are interested in computing $\langle HHLL \rangle$ where, 'heavy' operators are defined by $h_i \sim 1/g^2$
- In the semi-classical limit, $g^2 \rightarrow 0$,

$$\{f(\tau), \tau\} = -g^2 \sum_{i} \left[\kappa_{1,i} \delta(\tau - \tau_{2i-1}) + \kappa_{1,i} \delta(\tau - \tau_{2i}) - \kappa_{2,i} \left(sgn(\tau - \tau_{2i-1}) - sgn(\tau - \tau_{2i}) \right) \right] + g^2 \kappa_3$$

can be
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integral

Ward Identities

[PN, Sonner, Vielma '19]

 Within Schwarzian theory, one can derive the following 'Ward' identity

$$\sum_{i=1}^{m} \left[h_i \left(\delta(x - x_{2i-1}) + \delta(x - x_{2i}) \right) + \left(sgn(x - x_{2i-1}) - sgn(x - x_{2i}) \right) \partial_{x_{2i}} \right] \prod_{j=1}^{n} \left\langle \mathcal{O}_j \left(x_{2j-1}, x_{2j} \right) \right\rangle$$

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• The solution for $f(\tau)$ can be found by solving the Hill's equation:

Fion:
$$p''(\tau) + \frac{1}{2}T(\tau) p(\tau) = 0$$
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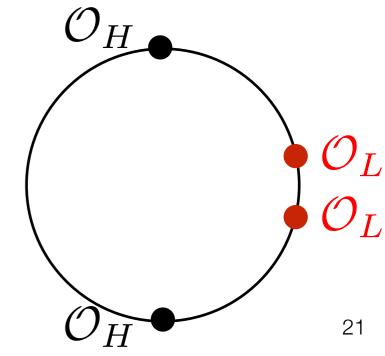
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but not when the background temperature goes to ZERO!

Effective temperature goes to ZERO with Background temperature

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[LMTV '18; Altland, Bagrets, Kamanev '16]

Schwarzian theory is closely related to 1-D
 Liouville theory, [Altland, Bagrets, Kamanev '16]
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Eigenfunctions of Liouville QM ≠ FZZ boundary states in 2D Liouville

These states *thermalize* with

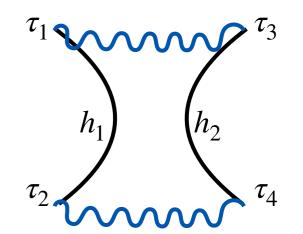
$$\tilde{T}_{eff} = \frac{\sqrt{2g^2 E}}{2\pi}$$

[Lam, Mertens, Turiaci, Verlinde '18]

ETH in the Conformal Limit

Correlation Functions

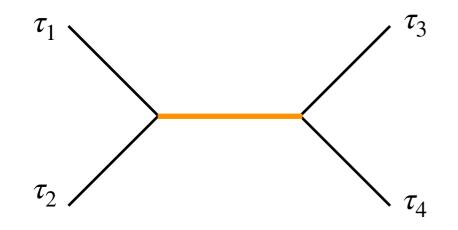
- Correlation functions of the fermions as well as of the excited operators in SYK spectrum get contribution from:
 - Exchange of Schwarzian modes

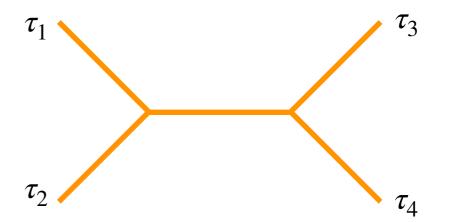


$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \mathcal{O}_4 \rangle =$$

$$\int \frac{f(\tau)}{\mathbb{SL}(2,\mathbb{R})} \exp \left[-\frac{1}{g^2} \int d\tau \left\{ f(\tau), \tau \right\} \right] \left(\frac{f'(\tau_1)f'(\tau_2)}{|f(\tau_1) - f(\tau_2)|^2} \right)^{h_1} \left(\frac{f'(\tau_3)f'(\tau_4)}{|f(\tau_3) - f(\tau_4)|^2} \right)^{h_2}$$

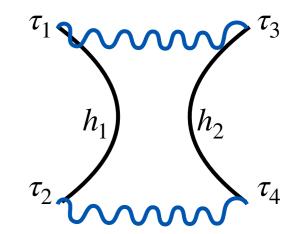
Exchange of excited modes (conformal limit)





Correlation Functions

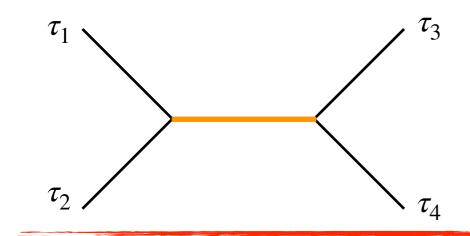
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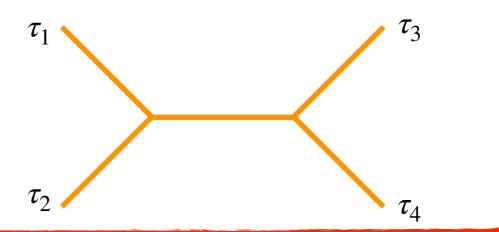


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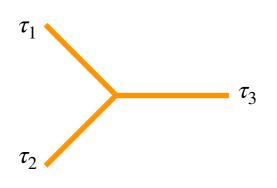
Exchange of excited modes (conformal limit)





A study of 3-point fins.

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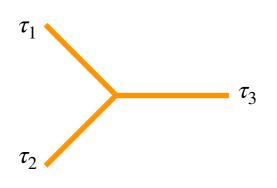


discrete tower of states
$$\mathcal{O}_n \sim \psi_i \; \partial^{2n+1} \psi_i$$

$$h_n = 2n+1+\epsilon_n$$

• In the conformal limit, we consider the 3-point correlation functions. [Jevicki, Suzuki, Yoon '16; Gross, Rosenhaus '17]

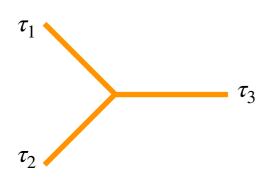
A study of 3-point fins.



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A Study of 3-point



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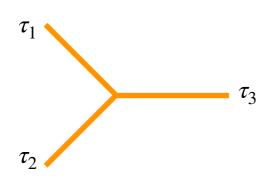
$$\langle \mathcal{O}_m \mathcal{O}_k \mathcal{O}_n \rangle = \mathfrak{g}(E, d) \times \left[2^{-2(2E-1)} \sqrt{2\pi E} \frac{\Gamma(4E-1)}{\Gamma(2E-d)\Gamma(2E+d)} \right]$$

$$\xrightarrow{E \to \infty} f_k(E) \, \delta_{m,n} + e^{-E \ln 2} \, R(E,d)$$

[PN, Sonner, Vielma '19]

$$E = \frac{m+n}{2}, \ d = n-m$$

A Study of 3-point fins.



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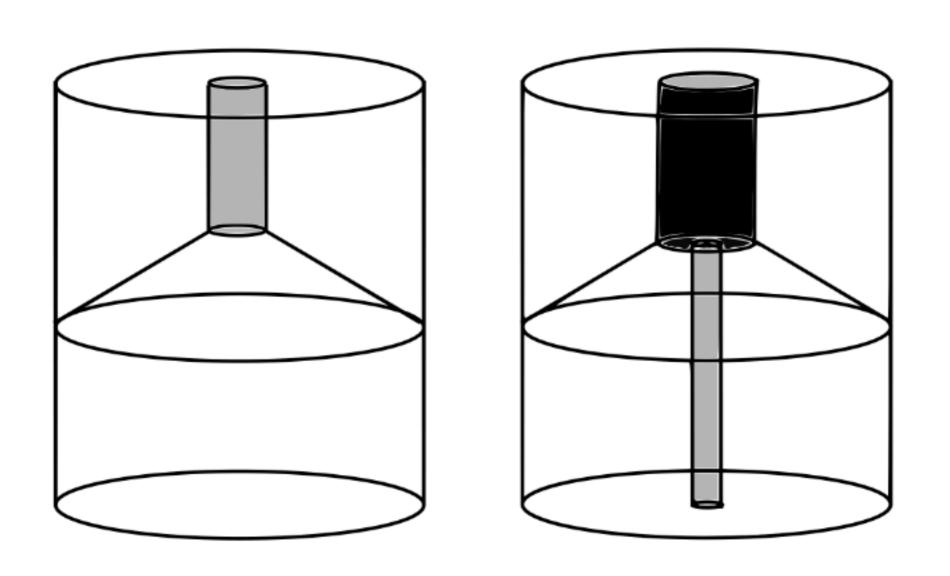
[PN, Sonner, Vielma '19]

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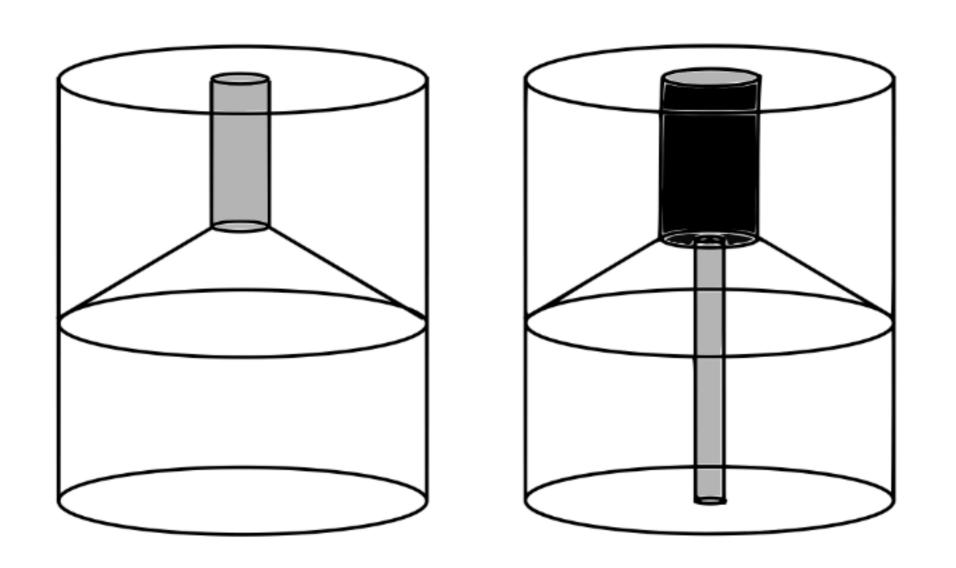
conformal sector of SYK satisfies ETH

A Bulle Story?

Dual of Schwarzian theory in 1st Limit?

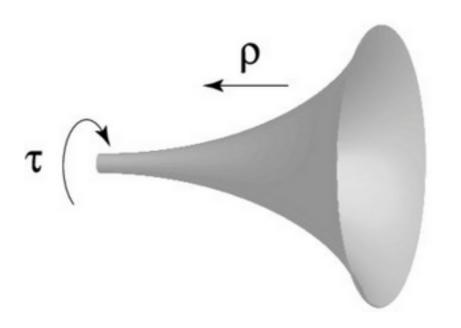


Dual of Schwarzian theory in 1st Limit?



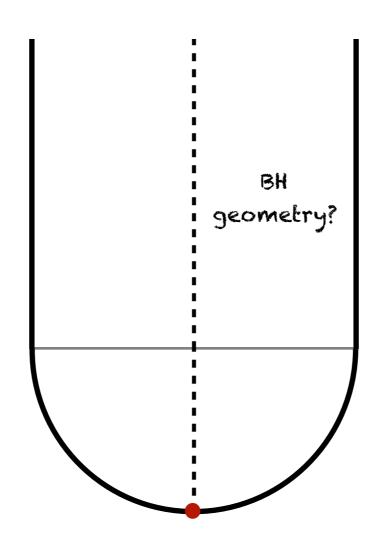
[Vos '18]

Dual of Schwarzian theory in 2nd Limit?

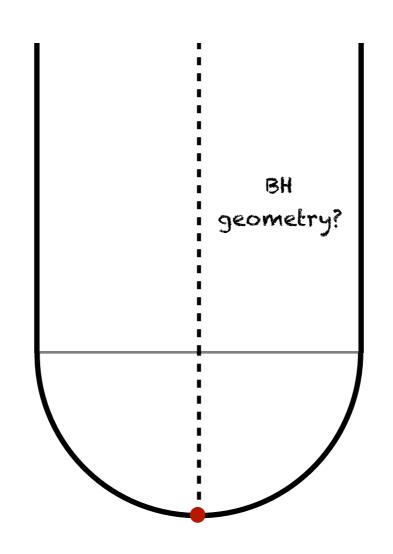


Thermal AdS₂

Dual of SYK theory in Conformal Limit?



Dual of SYK theory in Conformal Limit?



Can the ETH in conformal sector be understood as a correction?

Summary

- We showed that Eigenstate thermalization can be studied analytically in SYK model in 2 limits:
 - Schwarzian limit: where the contribution of pseudo-Goldstone modes is considered
 - First Thermodynamic limit: $N \to \infty$ then $\beta J \to \infty$ Weak ETH!
 - Second Thermodynamic limit: $N, \beta J \to \infty$ simultaneously, followed by $g^2 \sim \frac{\beta J}{N} \to 0$ NO ETH!
 - ETH in boundary states of Liouville theory: ETH!
 - ◆ Conformal limit: where only the operators that are primaries of Virasoro (and their descendants) are considered ETH!

TO DO

- Bulk dual of the each of the above limits
- Are 'heavy' states chaotic?

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